博士學位論文

# Eigenvalue estimates of the basic Dirac operator on a Riemannian foliation 

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# Eigenvalue estimates of the basic Dirac operator on a Riemannian foliation 

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# 엽층구조를 가지는 리만다양체에서의 

# basic Dirac 연산자의 고유치 계산 

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$<$ Abstract $>$

## Eigenvalue estimates of the basic Dirac operator on a Riemannian foliation

On a foliated Riemannian manifold with a transverse spin structure, we give a lower bound for the square of the eigenvalues of the basic Dirac operator in terms of the transversal scalar curvature and of the norm of an appropriate endomorphism of the normal bundle $Q$ of $\mathcal{F}$. We study the limiting case.

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## 1 Introduction

The Dirac operator on a Riemannian spin manifold, which were introduced by M.F. Atiyah and I.M. Singer([2]), were studied by many authors( $[3,4,7,9,10,15])$.

In 1963, A. Lichnerowicz([17]) proved that on a Riemannian spin manifold the square of the Dirac operator $D$ is given by

$$
D^{2}=\nabla^{*} \nabla+\frac{R}{4}
$$

where $\nabla^{*} \nabla$ is the positive spinor Laplacian and $R$ the scalar curvature. In particular, the first sharp estimate for the eigenvalues $\lambda$ of the Dirac operator $D$ was proved by Th. Friedrich ([7]) in 1980. Using a suitable Riemannian spin connection, he proved the inequality

$$
\begin{equation*}
\lambda^{2} \geq \frac{\text { 저저대하교 }}{4(n-1)} \inf _{M} \text { 주ㅇㅏㅏㄴ } \tag{1.1}
\end{equation*}
$$

on manifolds $\left(M^{n}, g\right)$ with positive scalar curvature $R>0$. He also proved, in the limiting case, that the manifold is an Einstein. The inequality (1.1) has been improved in several directions by many authors ([9,10,11,12]).

In 1988, J.Brüning and F.W.Kamber([5]) introduced the transversal spin structure on the normal bundle $Q$ of the foliation $\mathcal{F}$. Let $S(\mathcal{F})$ be a foliated spinor bundle on $(M, \mathcal{F})$. Then the transversal Dirac operator $D_{t r}: S(\mathcal{F}) \rightarrow S(\mathcal{F})$ is defined by

$$
\begin{equation*}
D_{t r} \Psi=\sum_{a} E_{a} \cdot \nabla_{E_{a}} \Psi-\frac{1}{2} \kappa \cdot \Psi \tag{1.2}
\end{equation*}
$$

where $\kappa$ is a mean curvature form of $\mathcal{F}$. If $\mathcal{F}$ is isoparametric and $\kappa$ is divergence-free, i.e. $\delta \kappa=0$, then the Lichnerowicz type formula is given
by

$$
\begin{equation*}
D_{t r}^{2} \Psi=\nabla_{t r}^{*} \nabla_{t r} \Psi+\frac{1}{4} K^{\sigma} \tag{1.3}
\end{equation*}
$$

where $K^{\sigma}=\sigma^{\nabla}+|\kappa|^{2}, \sigma^{\nabla}$ is the transversal scalar curvature of $\mathcal{F}$.
In 2001, Jung([11]) studied the eigenvalue of the basic Dirac operator and the limiting foliation. Namely, let $\left(M, g_{M}, \mathcal{F}\right)$ be a Riemannian manifold with an isoparametric transverse spin folition $\mathcal{F}$ of codimension $q>1$ and bundle-like metric $g_{M}$ with respect to $\mathcal{F}$. Assume that the mean curvature $\kappa$ of $\mathcal{F}$ satisfies $\delta \kappa=0$ and $K^{\sigma} \geq 0$. Then the eigenvalue $\lambda$ of the basic Dirac operator $D_{b}$ satisfies

$$
\begin{equation*}
\lambda^{2} \geq \frac{1}{4} \frac{q}{q-1} \inf _{M} K^{\sigma} . \tag{1.4}
\end{equation*}
$$

In the limiting case, it is proved that $\mathcal{F}$ is minimal, transversally Einsteinian with constant transversal scalar curvature.

In 2004, Jung et. al([12]) gave new lower bound for the eigenvalues of $D_{b}$ by the first eigenvalue of the basic Yamabe operator $Y_{b}$, which is defined by

$$
\begin{equation*}
Y_{b}=4 \frac{q-1}{q-2} \Delta_{B}+\sigma^{\nabla} \tag{1.5}
\end{equation*}
$$

where $\Delta_{B}$ is a basic Laplacian acting on basic functions. Namely, let $\left(M, g_{M}, \mathcal{F}\right)$ be a compact Riemannian manifold with a transverse spin foliation $\mathcal{F}$ of codimension $q \geq 3$ and bundle-like metric $g_{M}$ such that $\kappa \in \Omega_{B}^{1}(\mathcal{F})$ and $\delta \kappa=0$. Then

$$
\begin{equation*}
\lambda^{2} \geq \frac{q}{4(q-1)}\left(\mu_{1}+\inf |\kappa|^{2}\right), \tag{1.6}
\end{equation*}
$$

where $\mu_{1}$ is the first eigenvalue of $Y_{b}$.

In this thesis, we give a lower bound for the square of the eigenvalues of the basic Dirac operator in terms of the transversal scalar curvature and of the norm of an appropriate endomorphism of the normal bundle $Q$ of $\mathcal{F}$.

This article is organized as followings. In Chapter 2, we review the known fact on the foliated Riemannian manifold. In Chapter 3, we study some basic properties of the transversal Dirac operator $D_{t r}$. In Chapter 4, we estimate the conformal lower bound for the eigenvalues of the basic Dirac operator by the modified new connection. We apply some techniques and concerning conformal changes of the Riemannian metric to get a sharper estimate than the theorem in terms of the first eigenvalues of the Yamabe operator. Namely,

Theorem 1.1 $\operatorname{Let}\left(M, \tilde{g}_{M}, \mathcal{F}\right)$ be a compact Riemannian manifold with a transverse spin foliation $\mathcal{F}$ of codimension $q \geq 3$ and a bundle-like metric $\tilde{g}_{M}$. Then, any eigenvalue $\lambda$ of the basic Dirac operator corresponding to the eigenspinor $\Psi \in \Gamma S(\mathcal{F})$ satisfies

$$
\begin{equation*}
\lambda^{2} \geq \frac{1}{4}\left(\mu_{1}+\inf _{M}|\kappa|^{2}\right)+\inf _{M}\left|l_{\Psi}\right|^{2} \tag{1.7}
\end{equation*}
$$

where $\mu_{1}$ is the first eigenvalue of the basic Yamabe operator $Y_{b}$ of $\mathcal{F}$ and $l_{\Psi}$ is a symmetric endomorphism associated with $F_{\Psi}$.

In Chapter 5, we prove, in the limiting case, that the foliation $\mathcal{F}$ is minimal.

Theorem 1.2 Let $\left(M, \tilde{g}_{M}, \mathcal{F}\right)$ be a compact Riemannian manifold with a transverse spin foliation $\mathcal{F}$ of codimension $q \geq 3$ and a bundle-like metric
$\tilde{g}_{M}$. Assume that an eigenvalue $\lambda$ of $D_{b}$ corresponding to the eigenspinor $\Psi$ satisfies

$$
\lambda^{2}=\frac{1}{4}\left(\mu_{1}+\inf _{M}|\kappa|^{2}\right)+\inf _{M}\left|l_{\Psi}\right|^{2} .
$$

Then $\left|l_{\Psi}\right|$ is constant and the foliation $\mathcal{F}$ is minimal. Moreover

$$
\begin{equation*}
\left(\operatorname{div}_{\nabla} l_{\Psi}\right)(X)=(1-q) g_{Q}\left(l_{\Psi}(X), \operatorname{grad}_{\nabla}(u)\right) \tag{1.8}
\end{equation*}
$$

for any $X \in \Gamma Q$.

Throughout this paper, we consider the bundle-like metric $\tilde{g}_{M}$ for $(M, \mathcal{F})$ such that the mean curvature form of $\mathcal{F}$ is basic-harmonic. The existence of the bundle-like metric $g_{M}$ for $(M, \mathcal{F})$ such that $\kappa$ is basic, i.e. $\kappa \in$ $\Omega_{B}^{1}(\mathcal{F})$, is proved in [6]. In [18,19], it is proved that for any bundle-like metric $g_{M}$ with $\kappa \in \Omega_{B}^{1}(\mathcal{F})$ there exists another bundle-like metric $\tilde{g}_{M}$ for which the mean curvature form is basic-harmonic.

## 2 Riemannian foliation

### 2.1 Definition of foliations

Let $M$ be a smooth manifold of dimension $p+q$.
Definition 2.1 A codimension $q$ foliation $\mathcal{F}$ on $M$ is given by an open cover $\mathcal{U}=\left(U_{i}\right)_{i \in I}$ and for each $i$, a diffeomorphism $\varphi_{i}: \mathbb{R}^{p+q} \rightarrow U_{i}$ such that, on $U_{i} \cap U_{j} \neq \emptyset$, the coordinate change $\varphi_{j}^{-1} \circ \varphi_{i}: \varphi_{i}^{-1}\left(U_{i} \cap U_{j}\right) \rightarrow$ $\varphi_{j}^{-1}\left(U_{i} \cap U_{j}\right)$ has the form

$$
\begin{equation*}
\varphi_{j}^{-1} \circ \varphi_{i}(x, y)=\left(\varphi_{i j}(x, y), \gamma_{i j}(y)\right) \tag{2.1}
\end{equation*}
$$

Equivalently, we have the following another definition. Let $f_{i}=p r \circ \varphi_{i}^{-1}$ : $U_{i} \rightarrow \mathbb{R}^{q}$ be a submersion, where $p r: \mathbb{R}^{p+q} \rightarrow \mathbb{R}^{q}$ is a projection.

Definition 2.2 A codimension $q$ foliation $\mathcal{F}$ on $M$ is given by an open cover $\mathcal{U}=\left(U_{i}\right)_{i \in I}$, submersions $f_{i}: U_{i} \rightarrow N$ over $q$ dimensional model manifold $N$ and for $U_{i} \cap U_{j} \neq \emptyset$, a diffeomorphism (transition function) $\gamma_{i j}: f_{i}\left(U_{i} \cap U_{j}\right) \rightarrow f_{j}\left(U_{i} \cap U_{j}\right)$ satisfying

$$
\begin{equation*}
f_{j}(x)=\gamma_{i j} \circ f_{i}(x) \quad \text { for } x \in U_{i} \cap U_{j} . \tag{2.2}
\end{equation*}
$$

From Definition 2.1(or 2.2), the manifold $M$ is decomposed into connected submanifolds of dimension $p$. Each of these submanifolds is called a leaf of $\mathcal{F}$. Coordinate patches $\left(U_{i}, \varphi_{i}\right)$ are said to be distinguished for the foliation $\mathcal{F}$. The tangent bundle $L$ of a foliation is the subbundle of $T M$, consisting of all vectors tangent to the leaves of $\mathcal{F}$. The normal
bundle $Q$ of a codimension $q$ foliation $\mathcal{F}$ on $M$ is the quotient bundle $Q=T M / L$. Equivalently $Q$ appears in the exact sequence of vector bundles

$$
\begin{equation*}
0 \rightarrow L \rightarrow T M \xrightarrow{\pi} Q \rightarrow 0 \tag{2.3}
\end{equation*}
$$

If $\left(x_{1}, \ldots, x_{p} ; y_{1}, \ldots, y_{q}\right)$ are local coordinates in a distinguished chart $U$, the bundle $Q \mid U$ is framed by the vector fields $\pi \frac{\partial}{\partial y_{1}}, \ldots, \pi \frac{\partial}{\partial y_{q}}$. For a vector field $Y \in \Gamma T M$, we denote also $\bar{Y}=\pi Y \in \Gamma Q$. A vector field $Y$ on $U$ is projectable, if $Y=\sum_{i} a_{i} \frac{\partial}{\partial x_{i}}+\sum_{\alpha} b_{\alpha} \frac{\partial}{\partial y_{\alpha}}$ with $\frac{\partial b_{\alpha}}{\partial x_{i}}=0$ for all $\alpha=1, \ldots, q$ and $i=1, \ldots, p$. This means that the functions $b_{\alpha}=b_{\alpha}(y)$ are independent of $x$. Then $\bar{Y}=\sum_{\alpha} b_{\alpha} \frac{\bar{\partial}}{\partial y_{\alpha}}$ with $b_{\alpha}$ independent of $x$. This property is preserved under change of distinguished charts, hence makes intrinsic sense.

The transversal geometry of a foliation is the geometry infinitesimally modeled by $Q$, while the tangential geometry is infinitesimally modeled by $L$. A key fact is the existence of the Bott connection in $Q$ defined by

$$
\begin{equation*}
\stackrel{\circ}{\nabla}_{X} s=\pi\left(\left[X, Y_{s}\right]\right) \quad \text { for } X \in \Gamma L \tag{2.4}
\end{equation*}
$$

where $Y_{s} \in T M$ is any vector field projecting to $s$ under $\pi: T M \rightarrow$ $Q$. It is a partial connection along $L$. The right hand side in (2.4) is independent of the choice of $Y_{s}$. Namely, the difference of two such choices is a vector field $X^{\prime} \in \Gamma L$ and $\left[X, X^{\prime}\right] \in \Gamma L$ so that $\pi\left[X, X^{\prime}\right]=0$.

A Riemannian metric $g_{Q}$ on the normal bundle $Q$ of a foliation $\mathcal{F}$ is holonomy invariant, if

$$
\begin{equation*}
\theta(X) g_{Q}=0 \quad \text { for all } X \in \Gamma L \tag{2.5}
\end{equation*}
$$

where $\theta(X)$ is Lie derivative. Here we have by definition for $s, t \in \Gamma Q$,

$$
\left(\theta(X) g_{Q}\right)(s, t)=X g_{Q}(s, t)-g_{Q}(\theta(X) s, t)-g_{Q}(s, \theta(X) t)
$$

Definition 2.3 A Riemannian foliation is a foliation $\mathcal{F}$ with a holonomy invariant transversal metric $g_{Q}$. A metric $g_{M}$ is a bundle-like, if the induced metric $g_{Q}$ on $Q$ is holonomy invariant.

The study of Riemannian foliations was initiated by Reinhart in 1959([21]). A simple example of a Riemannian foliation is given by a nonsingular Killing vector field $X$ on $\left(M, g_{M}\right)$. This means that $\theta(X) g_{M}=0$.

An adapted connection in $Q$ is a connection restricting along $L$ to the partial Bott connection $\stackrel{\circ}{\nabla}$. To show that such connections exist, consider a Riemannian metric $g_{M}$ on $M$. Then $T M$-splits orthogonally as $T M=L \oplus L^{\perp}$. This means that there is a bundle map $\sigma: Q \rightarrow L^{\perp}$ splitting the exact sequence (2.3), i.e., satisfying $\pi \circ \sigma=$ identity. This metric $g_{M}$ on $T M$ is then a direct sum

$$
g_{M}=g_{L} \oplus g_{L^{\perp}}
$$

With $g_{Q}=\sigma^{*} g_{L^{\perp}}$, the splitting map $\sigma:\left(Q, g_{Q}\right) \rightarrow\left(L^{\perp}, g_{L^{\perp}}\right)$ is a metric isomorphism. Let now $\nabla^{M}$ be the Levi-Civita connection associated to the Riemannian metric $g_{M}$. Then the adapted connection $\nabla$ in $Q$ is defined by

$$
\left\{\begin{array}{l}
\nabla_{X} s=\stackrel{\circ}{\nabla}_{X} s=\pi\left(\left[X, Y_{s}\right]\right) \quad \text { for } X \in \Gamma L  \tag{2.6}\\
\nabla_{X} s=\pi\left(\nabla_{X}^{M} Y_{s}\right) \quad \text { for } X \in \Gamma L^{\perp}
\end{array}\right.
$$

where $s \in \Gamma Q$ and $Y_{s} \in \Gamma L^{\perp}$ corresponding to $s$ under the canonical isomorphism $Q \cong L^{\perp}$. For any connection $\nabla$ on $Q$, there is a torsion $T_{\nabla}$
defined by

$$
\begin{equation*}
T_{\nabla}(Y, Z)=\nabla_{Y} \pi(Z)-\nabla_{Z} \pi(Y)-\pi[Y, Z] \tag{2.7}
\end{equation*}
$$

for any $Y, Z \in \Gamma T M$. Then we have the following proposition ([22]).
Proposition 2.4 For any metric $g_{M}$ on $M$, and the adapted connection $\nabla$ on $Q$ defined by (2.6), we have $T_{\nabla}=0$.

Proof. For $X \in \Gamma L, Y \in \Gamma T M$ we have $\pi(X)=0$ and

$$
T_{\nabla}(X, Y)=\nabla_{X} \pi(Y)-\pi[X, Y]=0
$$

For $Z, Z^{\prime} \in \Gamma L^{\perp}$ we have

$$
T_{\nabla}\left(Z, Z^{\prime}\right)=\pi\left(\nabla_{Z}^{M} Z^{\prime}\right)-\pi\left(\nabla_{Z^{\prime}}^{M} Z\right)-\pi\left[Z, Z^{\prime}\right]=\pi\left(T_{\nabla^{M}}\left(Z, Z^{\prime}\right)\right)=0
$$

where $T_{\nabla^{M}}$ is the (vanishing) torsion of $\nabla^{M}$. Finally the bilinearity and skew symmetry of $T_{\nabla}$ imply the desired result.

The curvature $R^{\nabla}$ of $\nabla$ is defined by

$$
R^{\nabla}(X, Y)=\nabla_{X} \nabla_{Y}-\nabla_{Y} \nabla_{X}-\nabla_{[X, Y]} \quad \text { for } X, Y \in T M
$$

From an adapted connection $\nabla$ in $Q$ defined by (2.6), its curvature $R^{\nabla}$ coincides with $\stackrel{\circ}{R}$ for $X, Y \in \Gamma L$, hence $R^{\nabla}(X, Y)=0$ for $X, Y \in \Gamma L$. And we have the following proposition ([13,14,22]).

Proposition 2.5 Let $\left(M, g_{M}, \mathcal{F}\right)$ be a $(p+q)$-dimensional Riemannian manifold with a foliation $\mathcal{F}$ of codimension $q$ and bundle-like metric $g_{M}$ with respect to $\mathcal{F}$. Let $\nabla$ be a connection defined by (2.6) in $Q$ with curvature $R^{\nabla}$. Then for $X \in \Gamma L$ the following holds:

$$
\begin{equation*}
i(X) R^{\nabla}=\theta(X) R^{\nabla}=0 \tag{2.8}
\end{equation*}
$$

Proof. (i) Let $Y \in \Gamma T M$ and $s \in \Gamma Q$. Then

$$
\begin{aligned}
R^{\nabla}(X, Y) s & =\nabla_{X} \nabla_{Y} s-\nabla_{Y} \nabla_{X} s-\nabla_{[X, Y]} s \\
& =\theta(X) \nabla_{Y} s-\nabla_{Y} \theta(X) s-\nabla_{\theta(X) Y} s \\
& =(\theta(X) \nabla)_{Y} s=0 .
\end{aligned}
$$

(ii) Let $Y, Z \in \Gamma T M$ and $s \in \Gamma Q$. Then

$$
\begin{aligned}
\left(\theta(X) R^{\nabla}\right)(Y, Z) s= & \theta(X) R^{\nabla}(Y, Z) s-R^{\nabla}(\theta(X) Y, Z) s \\
& -R^{\nabla}(Y, \theta(X) Z) s-R^{\nabla}(Y, Z) \theta(X) s \\
= & \theta(X)\left\{\nabla_{Y} \nabla_{Z} s-\nabla_{Z} \nabla_{Y} s-\nabla_{[Y, Z]} s\right\} \\
& -\left\{\nabla_{\theta(X) Y} \nabla_{Z} s-\nabla_{Z} \nabla_{\theta(X) Y} s-\nabla_{[\theta(X) Y, Z]} s\right\} \\
- & \left\{\nabla_{Y} \nabla_{\theta(X) Z} s-\nabla_{\theta(X) Z} \nabla_{Y} s-\nabla_{[Y, \theta(X) Z]} s\right\} \\
& -\left\{\nabla_{Y} \nabla_{Z} \theta(X) s-\nabla_{Z} \nabla_{Y} \theta(X) s-\nabla_{[Y, Z]} \theta(X) s\right\} \\
= & \nabla_{Y}\left(\theta(X) \nabla_{Z} s\right)-\nabla_{Z}\left(\theta(X) \nabla_{Y} s\right)-\nabla_{\theta(X)[Y, Z]} s \\
& +\nabla_{Z} \nabla_{\theta(X) Y} s+\nabla_{[\theta(X) Y, Z]} s-\nabla_{Y} \nabla_{\theta(X) Z} s \\
& +\nabla_{[Y, \theta(X) Z]} s-\nabla_{Y} \nabla_{Z} \theta(X) s+\nabla_{Z} \nabla_{Y} \theta(X) s \\
= & -\nabla_{\theta(X)[Y, Z]} s+\nabla_{[\theta(X) Y, Z]} s+\nabla_{[Y, \theta(X) Z]} s \\
= & \left(-\nabla_{[X,[Y, Z]]}+\nabla_{[[X, Y], Z]}+\nabla_{[Y,[X, Z]]}\right) s=0 .
\end{aligned}
$$

By Proposition 2.5, we can define the (transversal) Ricci curvature $\rho^{\nabla}$ : $\Gamma Q \rightarrow \Gamma Q$ and the (transversal) scalar curvature $\sigma^{\nabla}$ of $\mathcal{F}$ by

$$
\begin{equation*}
\rho^{\nabla}(s)=\sum_{a} R^{\nabla}\left(s, E_{a}\right) E_{a}, \quad \sigma^{\nabla}=\sum_{a} g_{Q}\left(\rho^{\nabla}\left(E_{a}\right), E_{a}\right), \tag{2.9}
\end{equation*}
$$

where $\left\{E_{a}\right\}_{a=1, \cdots, q}$ is an orthonormal basis of $Q$.

Definition 2.6 The foliation $\mathcal{F}$ is said to be (transversally) Einsteinian if the model space $N$ is Einsteinian, that is,

$$
\begin{equation*}
\rho^{\nabla}=\frac{1}{q} \sigma^{\nabla} \cdot i d \tag{2.10}
\end{equation*}
$$

with constant transversal scalar curvature $\sigma^{\nabla}$.

### 2.2 Mean curvature form and basic Laplacian

The second fundamental form $\alpha$ of $\mathcal{F}$ is given by

$$
\begin{equation*}
\alpha(X, Y)=\pi\left(\nabla_{X}^{M} Y\right) \quad \text { for } X, Y \in \Gamma L \tag{2.11}
\end{equation*}
$$

Proposition $2.7 \alpha$ is $Q$-valued, bilinear and symmetric.
Proof. By definition, it is trivial that $\alpha$ is $Q$-valued and bilinear. Next, by torsion freeness of $\nabla^{M}$, we have that for any $X, Y \in \Gamma L$,

$$
\alpha(X, Y)=\pi\left(\nabla_{X}^{M} Y\right)=\pi\left(\nabla_{Y}^{M} X\right)-\pi([X, Y])
$$

Since $[X, Y] \in \Gamma L$ for any $X, Y \in \Gamma L$, we have

$$
\alpha(X, Y)=\pi\left(\nabla_{Y}^{M} X\right)=\alpha(Y, X)
$$

Definition 2.8 The mean curvature vector field of $\mathcal{F}$ is then defined by

$$
\begin{equation*}
\tau=\sum_{i} \alpha\left(E_{i}, E_{i}\right)=\sum_{i} \pi\left(\nabla_{E_{i}}^{M} E_{i}\right) \tag{2.12}
\end{equation*}
$$

where $\left\{E_{i}\right\}_{i=1, \cdots, p}$ is an orthonormal basis of $L$. The dual form $\kappa$, the mean curvature form for $L$, is then given by

$$
\begin{equation*}
\kappa(X)=g_{Q}(\tau, X) \quad \text { for } X \in \Gamma Q \tag{2.13}
\end{equation*}
$$

The foliation $\mathcal{F}$ is said to be minimal (or harmonic ) if $\kappa=0$.

Definition 2.9 Let $\mathcal{F}$ be an arbitrary foliation on a manifold $M$. A differential form $\omega \in \Omega^{r}(M)$ is basic, if

$$
\begin{equation*}
i(X) \omega=0, \theta(X) \omega=0, \text { for } X \in \Gamma L \tag{2.14}
\end{equation*}
$$

In a distinguished chart $\left(x_{1}, \ldots, x_{p} ; y_{1}, \ldots, y_{q}\right)$ of $\mathcal{F}$, a basic form $w$ is expressed by

$$
\omega=\sum_{a_{1}<\cdots<a_{r}} \omega_{a_{1} \cdots a_{r}} d y_{a_{1}} \wedge \cdots \wedge d y_{a_{r}},
$$

where the functions $\omega_{a_{1} \cdots a_{r}}$ are independent of $x$, i.e. $\frac{\partial}{\partial x_{i}} \omega_{a_{1} \cdots a_{r}}=0$. Let $\Omega_{B}^{r}(\mathcal{F})$ be the set of all basic $r$-forms on $M$. The exterior derivative $d$ preserves basic forms, since $\theta(X) d \omega=d \theta(X) \omega=0, i(X) d \omega=\theta(X) \omega-$ $d i(X) \omega=0$ for a basic form $\omega$. Hence $\Omega_{B}^{r}(\mathcal{F})$ constitutes a subcomplex

$$
d: \Omega_{B}^{r}(\mathcal{F}) \rightrightarrows \Omega_{B}^{r+1}(\mathcal{F}) \text { 도서관 }
$$

of the De Rham complex $\Omega^{*}(M)$ and the restriction $d_{B}=\left.d\right|_{\Omega_{B}^{*}(\mathcal{F})}$ is well defined. Its cohomology

$$
H_{B}(\mathcal{F})=H\left(\Omega_{B}^{*}(\mathcal{F}), d_{B}\right)
$$

is the basic cohomology of $\mathcal{F}$. It plays the role of the De Rham cohomology of the leaf space $M / \mathcal{F}$ of the foliation. Let $\delta_{B}$ the formal adjoint operator of $d_{B}$. Then we have the following proposition $([1,11])$.

Proposition 2.10 On a Riemannian foliation $\mathcal{F}$, we have

$$
\begin{equation*}
d_{B}=\sum_{a} \theta_{a} \wedge \nabla_{E_{a}}, \quad \delta_{B}=-\sum_{a} i\left(E_{a}\right) \nabla_{E_{a}}+i\left(\kappa_{B}\right), \tag{2.15}
\end{equation*}
$$

where $\kappa_{B}$ is the basic component of $\kappa,\left\{E_{a}\right\}$ is a local orthonormal basic frame in $Q$ and $\left\{\theta_{a}\right\}$ its $g_{Q}$-dual 1-form.

The foliation $\mathcal{F}$ is said to be isoparametric if $\kappa \in \Omega_{B}^{1}(\mathcal{F})$. We already know that $\kappa$ is closed, i.e., $d \kappa=0$ if $\mathcal{F}$ is isoparametric ([18]).

Definition 2.11 The basic Laplacian acting on $\Omega_{B}^{*}(\mathcal{F})$ is defined by

$$
\begin{equation*}
\Delta_{B}=d_{B} \delta_{B}+\delta_{B} d_{B} \tag{2.16}
\end{equation*}
$$

The following theorem is proved in the same way as the corresponding usual result in De Rham-Hodge Theory.

Theorem 2.12 ([22]) Let $\mathcal{F}$ be a transversally oriented Riemannian foliation on a closed oriented manifold $\left(M, g_{M}\right)$. Assume $g_{M}$ to be bundlelike metric with $\kappa \in \Omega_{B}^{1}(\mathcal{F})$. Then

$$
H_{B}^{r}(\mathcal{F}) \cong \mathcal{H}_{B}^{r}(\mathcal{F})
$$

where $\mathcal{H}_{B}^{r}(\mathcal{F})=\left\{\omega \in \Omega^{r}(\bar{M}) \mid \bar{\Delta}_{B}[\omega \mid\right.$ 희 0$\}$. 중앙도서관
If $\mathcal{F}$ is the foliation by points of $M$, the basic Laplacian is the ordinary Laplacian. In the more general case, the basic Laplacian and its spectrum provide information about the transverse geometry of $(M, \mathcal{F})([16])$.

### 2.3 Transversal divergence theorem

For the later use, we recall the divergence theorem on a foliated Riemannian manifold ([23]).

Theorem 2.13 $\operatorname{Let}\left(M, g_{M}, \mathcal{F}\right)$ be a closed, oriented, connected Riemannian manifold with a transversally orientable foliation $\mathcal{F}$ and a bundlelike metric $g_{M}$ with respect to $\mathcal{F}$. Then

$$
\begin{equation*}
\int_{M} d i v_{\nabla}(X)=\int_{M} g_{Q}(X, \tau) \tag{2.17}
\end{equation*}
$$

for all $X \in \Gamma Q$, where $\operatorname{div}_{\nabla}(X)$ denotes the transverse divergence of $X$ with respect to the connection $\nabla$ defined by (2.6).

Proof. Let $\left\{E_{i}\right\}$ and $\left\{E_{a}\right\}$ be orthonormal basis of $L$ and $Q$, respectively. Then for any $X \in \Gamma Q$,

$$
\begin{aligned}
\operatorname{div}(X) & =g_{M}\left(\nabla_{E_{i}}^{M} X, E_{i}\right)+g_{M}\left(\nabla_{E_{a}}^{M} X, E_{a}\right) \\
& =-g_{M}\left(X, \pi\left(\nabla_{E_{i}}^{M} E_{i}\right)\right)+g_{M}\left(\pi\left(\nabla_{E_{a}}^{M} X\right), E_{a}\right) \\
& =-g_{Q}(X, \tau)+g_{Q}\left(\nabla_{E_{a}} X, E_{a}\right) \\
& =-g_{Q}(X, \tau)+\operatorname{div}_{\nabla}(X) .
\end{aligned}
$$

By Green's Theorem on an ordinary manifold $M$, we have

$$
0=\int_{M} \operatorname{div}(X) d_{M}=\int_{M} \operatorname{div}_{M}(X) d_{M-}-\int_{M} g_{Q}(X, \tau) .
$$

This completes the proof of this Theorem.

Corollary 2.14 If $\mathcal{F}$ is minimal, then we have that for any $X \in \Gamma Q$,

$$
\begin{equation*}
\int_{M} d i v_{\nabla}(X)=0 \tag{2.18}
\end{equation*}
$$

## 3 Transverse spin structure

### 3.1 Clifford algebras

Definition 3.1 Let $V$ be a vector space over a field $K=\{\mathbb{R}, \mathbb{C}\}$ of dimension $n$ and $g$ a non-degenerate bilinear form on $V$. The Clifford algebra $C l(V, g)$ associated to $g$ on $V$ is the algebra over $K$ generated by $V$ with the relation

$$
\begin{equation*}
v \cdot w+w \cdot v=-2 g(v, w) 1 \tag{3.1}
\end{equation*}
$$

for $v, w \in V$. The product "." is called the Clifford multiplication.

Equivalently, the Clifford algebra of $V$ is given by the following universal problem(for detail, see [15])체주대학교 중앙도서관

Proposition 3.2 (Universal property) Let $A$ be an associative algebra with unit and $f: V \rightarrow A$ a linear map such that for all $v \in V$

$$
f(v)^{2}=-g(v, v) 1
$$

Then $f$ uniquely extends to a $K$-algebra homomorphism

$$
\tilde{f}: C l(V, g) \rightarrow A
$$

Remark. The Clifford algebra may be realized as the quotient

$$
C l(V, g):=T(V) / I(V, g)
$$

where $T(V)$ is the tensor algebra of $V$, and $I(V, g)$ the ideal generated by all elements of the form $v \otimes v+g(v, v) 1$, for $v \in V$.

Remark. (1) If $\left(E_{1}, \cdots E_{n}\right)$ is a g-orthonormal basis of $V$, then

$$
\left\{E_{i_{1}} \cdot \ldots \cdot E_{i_{k}} \mid 1 \leq i_{1}<\ldots<i_{k} \leq n, \quad 0 \leq k \leq n\right\}
$$

is a basis of $C l(V, g)$, thus $\operatorname{dim} C l(V, g)=2^{n}$.
(2) There is a canonical isomorphism of vector spaces, between the exterior algebra and the Clifford algebra of $(V, g)$ which is given by :

$$
\begin{gathered}
\wedge^{*} V \stackrel{\sim}{\rightrightarrows} C l(V, g) \\
E_{i_{1}} \wedge \ldots \wedge E_{i_{k}} \longmapsto E_{i_{1}} \cdot \ldots \cdot E_{i_{k}} .
\end{gathered}
$$

This isomorphism does not depend on the choice of the basis. Let us denote $C l_{n}=C l\left(\mathbb{R}^{n},<,>\right)$. Then we have the following proposition ([15]).

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Proposition 3.3 For all $v \in \mathbb{R}^{n}$ and all $\varphi \in C l_{n}$, we have

$$
v \cdot \varphi \simeq v \wedge \varphi-i(v) \varphi \quad \text { and } \quad \varphi \cdot v \simeq(-1)^{p}(v \wedge \varphi+i(v) \varphi),
$$

where $\wedge$ denotes the exterior, $i(v)$ the interior product and $\varphi \in \wedge^{p} \mathbb{R}^{n} \subset$ $\wedge^{*} \mathbb{R}_{n} \simeq C l_{n}$.

Proof. Let $v=E_{j}$ and $\varphi=E_{i_{1}} \cdot \ldots \cdot E_{i_{p}}$.

1. If there exists $i_{k}$ such that $j=i_{k}$ then $v \wedge \varphi=0$ and

$$
\begin{aligned}
i(v) \varphi & =(-1)^{k-1} E_{i_{1}} \wedge \ldots \wedge E_{i_{k-1}} \wedge E_{i_{k+1}} \wedge \ldots E_{i_{p}} \\
& \simeq(-1)^{k-1} E_{i_{1}} \cdot \ldots \cdot E_{i_{k-1}} \cdot E_{i_{k+1}} \cdot \ldots \cdot E_{i_{p}} \\
& =-v \cdot \varphi \\
& =(-1)^{p} \varphi \cdot v .
\end{aligned}
$$

2. If $j \notin\left\{i_{1}, \ldots, i_{p}\right\}$ then $i(v) \varphi=0$ and

$$
\begin{aligned}
v \wedge \varphi & =E_{j} \wedge E_{i_{1}} \wedge \ldots \wedge E_{i_{p}} \simeq E_{j} \cdot E_{i_{1}} \cdot \ldots \cdot E_{i_{p}} \\
& =v \cdot \varphi \\
& =(-1)^{p} \varphi \cdot v .
\end{aligned}
$$

As the equalities of the assertion are bilinear, the proposition is proved.

Definition 3.4 The Pin group $\operatorname{Pin}(V)$ is defined by

$$
\begin{equation*}
\operatorname{Pin}(V)=\left\{a \in C l(V) \mid a=a_{1} \cdots a_{k},\left\|a_{i}\right\|=1\right\} . \tag{3.2}
\end{equation*}
$$

The Spin group is defined by

$$
\begin{gather*}
\text { 제주대학교 중앙도서관 }  \tag{3.3}\\
\operatorname{Spin}(V)=\left\{a \in \operatorname{Pin}(V) \mid a a^{t}=1\right\},
\end{gather*}
$$

where $a^{t}=a_{k} \cdots a_{1}$ for any $a=a_{1} \cdots a_{k}$. Equivalently, $\operatorname{Spin}(V)=$ $\left\{e_{1} \cdots e_{2 k}|\quad| e_{i} \mid=1\right\}$.

Let $V$ be a real vector space. Then $\operatorname{Spin}(V)$ is a compact and connected Lie group, and for $\operatorname{dim} V \geq 3$, it is also simply connected. Thus, for $\operatorname{dim} V \geq 3, \operatorname{Spin}(V)$ is the universal cover of $S O(V)$ (for detail, see [15]).

### 3.2 Transversal Dirac operator

Let $\left(M, g_{M}, \mathcal{F}\right)$ be a Riemannian manifold with a transversally oriented Riemannian foliation $\mathcal{F}$ of codimension $q$ and a bundle-like metric $g_{M}$ with respect to $\mathcal{F}$. Let $S O(q) \rightarrow P_{S O} \rightarrow M$ be the principal
bundle of (oriented) transverse orthonormal framings. Then a transverse spin structure is a principal $\operatorname{Spin}(q)$-bundle $P_{\text {Spin }}$ together with two sheeted covering $\xi: P_{\text {Spin }} \rightarrow P_{S O}$ such that $\xi(p \cdot g)=\xi(p) \xi_{0}(g)$ for all $p \in P_{\text {Spin }}, g \in \operatorname{Spin}(q)$, where $\xi_{0}: \operatorname{Spin}(q) \rightarrow S O(q)$ is a covering. In this case, the foliation $\mathcal{F}$ is called a transverse spin foliation. We then define the foliated spinor bundle $S(\mathcal{F})$ associated with $P_{\text {Spin }}$ by

$$
\begin{equation*}
S(\mathcal{F})=P_{\text {Spin }} \times \times_{\text {Spin }(q)} S_{q}, \tag{3.4}
\end{equation*}
$$

where $S_{q}$ is the irreducible spinor space associated to $Q$. The Hermitian metric $<,>$ on $S(\mathcal{F})$ induced from $g_{Q}$ satisfies the following relation:

$$
\begin{equation*}
<\varphi, \psi>=<v \cdot \varphi, v \cdot \psi> \tag{3.5}
\end{equation*}
$$

for every $v \in Q, g_{Q}(v, v)$ TH and $\bar{\varphi}, \psi \in S_{q-}$ And the Riemannian connection $\nabla$ on $P_{S O}$ defined by (2.6) can be lifted to one on $P_{S p i n}$, in particular, to one on $S(\mathcal{F})$, which will be denoted by the same letter.

Proposition 3.5 ([11, 15]) The spinorial covariant derivative on $S(\mathcal{F})$ is given locally by:

$$
\begin{equation*}
\nabla \Psi_{\alpha}=\frac{1}{4} \sum_{a, b} g_{Q}\left(\nabla E_{a}, E_{b}\right) E_{a} \cdot E_{b} \cdot \Psi_{\alpha} \tag{3.6}
\end{equation*}
$$

where $\Psi_{\alpha}$ is an orthonormal basis of $S_{q}$. And the curvature transform $R^{S}$ on $S(\mathcal{F})$ is given as

$$
\begin{equation*}
R^{S}(X, Y) \Phi=\frac{1}{4} \sum_{a, b} g_{Q}\left(R^{\nabla}(X, Y) E_{a}, E_{b}\right) E_{a} \cdot E_{b} \cdot \Phi \quad \text { for } X, Y \in T M \tag{3.7}
\end{equation*}
$$

where $\left\{E_{a}\right\}$ is an orthonormal basis of the normal bundle $Q$.

Proposition 3.6 ([15]) (Compatibility of $\nabla$ with " $\cdot$ " and $<\cdot, \cdot>$ )
(1) $X<\psi, \varphi>=<\nabla_{X} \psi, \varphi>+<\psi, \nabla_{X} \varphi>, \quad X \in \Gamma T M$
(2) $\quad \nabla_{X}(Y \cdot \psi)=\left(\nabla_{X} Y\right) \cdot \psi+Y \cdot \nabla_{X} \psi, \quad Y \in \Gamma Q$.

Theorem 3.7 ([11,12]) On the foliated spinor bundle $S(\mathcal{F})$, we have

$$
\begin{align*}
& \sum_{a} E_{a} \cdot R^{S}\left(X, E_{a}\right) \Phi=-\frac{1}{2} \rho^{\nabla}(\pi(X)) \cdot \Phi,  \tag{3.10}\\
& \sum_{a<b} E_{a} \cdot E_{b} \cdot R^{S}\left(E_{a}, E_{b}\right) \Phi=\frac{1}{4} \sigma^{\nabla} \Phi \tag{3.11}
\end{align*}
$$

for $X \in T M$.
Taking $\hat{\pi}$ to denote the projection

$$
\hat{\pi}: C^{\infty}\left(T^{*} M \otimes S(\mathcal{F})\right) \rightarrow C^{\infty}\left(Q^{*} \otimes S(\mathcal{F})\right) \cong C^{\infty}(Q \otimes S(\mathcal{F}))
$$

we define the transversal Dirac Operator $D_{t r}^{\prime}([5,8])$ by

$$
D_{t r}^{\prime}=\cdot \circ \hat{\pi} \circ \nabla
$$

If $\left\{E_{a}\right\}_{a=1, \cdots, q}$ is taken to be a local orthonormal basic frame in $Q$, then

$$
D_{t r}^{\prime}=\sum_{a} E_{a} \cdot \nabla_{E_{a}} .
$$

In $[5,8]$ it was shown that the formal adjoint $D_{t r}^{\prime *}$ is given by $D_{t r}^{\prime *}=$ $D_{t r}^{\prime}-\kappa \cdot$ and that therefore

$$
\begin{equation*}
D_{t r}=D_{t r}^{\prime}-\frac{1}{2} \kappa . \tag{3.12}
\end{equation*}
$$

is a symmetric, transversally elliptic differential operator, with symbol $\sigma_{D_{t r}}$ satisfying $\sigma_{D_{t r}}(x, \xi)=\xi$ for $\xi \in Q_{x}^{*}$ and $\sigma_{D_{t r}}(x, \xi)=0$ for $\xi \in L_{x}^{*}$. Then the transversal Dirac operator $D_{t r}$ is locally defined by

$$
\begin{equation*}
D_{t r} \Psi=\sum_{a} E_{a} \cdot \nabla_{E_{a}} \Psi-\frac{1}{2} \kappa \cdot \Psi \quad \text { for } \Psi \in \Gamma S(\mathcal{F}), \tag{3.13}
\end{equation*}
$$

where $\left\{E_{a}\right\}$ is a local orthonormal basic frame of $Q$. Then we have the Lichnerowicz-type formula on $\mathcal{F}$.

Theorem 3.8 On an isoparametric transverse spin foliation $\mathcal{F}$ with $\delta \kappa=$ 0 , it is well-known $([8,11])$ that

$$
\begin{equation*}
D_{t r}^{2} \Psi=\nabla_{t r}^{*} \nabla_{t r} \Psi+\frac{1}{4} K^{\sigma} \Psi \tag{3.14}
\end{equation*}
$$

where $K^{\sigma}=\sigma^{\nabla}+|\kappa|^{2}$ and

$$
\begin{equation*}
\nabla_{t r}^{*} \nabla_{t r} \Psi=-\sum_{a} \nabla_{E_{a}, E_{a}}^{2} \Psi+\nabla_{\kappa} \Psi . \tag{3.15}
\end{equation*}
$$

The operator $\nabla_{t r}^{*} \nabla_{t r}$ is non-negative and formally self-adjoint ([11]). In fact, we have the following proposition.

Proposition 3.9 ([11]) Let $\left(M, g_{M}, \mathcal{F}, S(\mathcal{F})\right)$ be a compact Riemannian manifold with the transverse spin foliation $\mathcal{F}$ and a bundle-like metric $g_{M}$ with respect to $\mathcal{F}$. Then

$$
\ll \nabla_{t r}^{*} \nabla_{t r} \Phi, \Psi \gg=\ll \nabla_{t r} \Phi, \nabla_{t r} \Psi \gg
$$

for all $\Phi, \Psi \in \Gamma E$, where $\ll \Phi, \Psi \gg=\int_{M}<\Phi, \Psi>$ is the inner product on $S(\mathcal{F})$.

We define the subspace $\Gamma_{B}(S(\mathcal{F}))$ of basic or holonomy invariant sections of $S(\mathcal{F})$ by

$$
\Gamma_{B}(S(\mathcal{F}))=\left\{\Psi \in \Gamma S(\mathcal{F}) \mid \nabla_{X} \Psi=0 \quad \text { for } X \in \Gamma L\right\}
$$

Trivially, we see that $D_{t r}$ leaves $\Gamma_{B}(S(\mathcal{F}))$ invariant if and only if the foliation $\mathcal{F}$ is isoparametric, i.e., $\kappa \in \Omega_{B}^{1}(\mathcal{F})$. Let $D_{b}=\left.D_{t r}\right|_{\Gamma_{B}(S(\mathcal{F}))}$ : $\Gamma_{B}(S(\mathcal{F})) \rightarrow \Gamma_{B}(S(\mathcal{F}))$. This operator $D_{b}$ is called the basic Dirac operator on (smooth) basic sections. It is well known that $D_{b}$ and $D_{b}^{2}$ have a discrete spectrum, respectively.

### 3.3 Transversal Dirac operators of transversally conformally related metrics

Now, we consider, for any real basic function $u$ on $M$, the transversally conformal metric $\bar{g}_{Q}=e^{2 u} g_{Q}$. Let $\bar{P}_{s o}(\mathcal{F})$ be the principal bundle of $\bar{g}_{Q}$ -orthogonal frames. Locally, the section $\bar{s}$ of $\bar{P}_{s o}(\mathcal{F})$ corresponding a section $s=\left(E_{1}, \cdots, E_{q}\right)$ of $P_{s o}(\mathcal{F})$ is $\bar{s}=\left(\bar{E}_{1}, \cdots, \bar{E}_{q}\right)$, where $\bar{E}_{a}=$ $e^{-u} E_{a}(a=1, \cdots, q)$. This isometry will be denoted by $I_{u}$. Thanks to the isomorphism $I_{u}$ one can define a transverse spin structure $\bar{P}_{\text {spin }}(\mathcal{F})$ on $\mathcal{F}$ in such a way that the diagram

commutes.
Let $\bar{S}(\mathcal{F})$ be the foliated spinor bundles associated with $\bar{P}_{\text {spin }}(\mathcal{F})$. For any section $\Psi$ of $S(\mathcal{F})$, we write $\bar{\Psi} \equiv I_{u} \Psi$. If $<,>_{g_{Q}}$ and $<,>_{\bar{g}_{Q}}$ denote respectively the natural Hermitian metrics on $S(\mathcal{F})$ and $\bar{S}(\mathcal{F})$, then for any $\Phi, \Psi \in \Gamma S(\mathcal{F})$

$$
\begin{equation*}
<\Phi, \Psi>_{g_{Q}}=<\bar{\Phi}, \bar{\Psi}>_{\bar{g}_{Q}} \tag{3.16}
\end{equation*}
$$

and the Clifford multiplication in $\bar{S}(\mathcal{F})$ is given by

$$
\begin{equation*}
\bar{X} \div \bar{\Psi}=\overline{X \cdot \Psi} \quad \text { for } X \in \Gamma Q \tag{3.17}
\end{equation*}
$$

Let $\bar{\nabla}$ be the metric and torsion free connection corresponding to $\bar{g}_{Q}$. Then we have the following proposition([12]).

Proposition 3.10 On the Riemannian foliation, we have that for $X, Y \in$ ГТ $M$,

$$
\begin{equation*}
\bar{\nabla}_{X} \pi(Y)=\nabla_{X} \pi(Y)+X(u) \pi(Y)+Y(u) \pi(X)-g_{Q}(\pi(X), \pi(Y)) \operatorname{grad}_{\nabla}(u) \tag{3.18}
\end{equation*}
$$

where $\operatorname{grad}_{\nabla}(u)=\sum_{a} E_{a}(u) E_{a}$ is a transversal gradient of $u$ and $X(u)$ is the Lie derivative of the function $u$ in the direction of $X$.

Proof. Since $\bar{\nabla}$ is the metric and torsion free connection with respect to $\bar{g}_{Q}$ on $Q$, we have

$$
\begin{aligned}
2 \bar{g}_{Q}\left(\bar{\nabla}_{X} s, t\right) & =X \bar{g}_{Q}(s, t)+Y \bar{g}_{Q}(\pi(X), t)-Z_{t} \bar{g}_{Q}(\pi(X), s) \\
& =\bar{g}_{Q}\left(\pi\left[X, Y_{s}\right], t\right)+\bar{g}_{Q}\left(\pi\left[Z_{t}, X\right], s\right)-\bar{g}_{Q}\left(\pi\left[Y_{s}, Z_{t}\right], \pi(X)\right)
\end{aligned}
$$

where $\pi\left(Y_{s}\right)=s$ and $\pi\left(Z_{t}\right)$ ㅔㅈㅏ $t$. From this formula, the proof is completed.

From (3.18), we have the following proposition([12]).
Proposition 3.11 The connection $\nabla$ and $\bar{\nabla}$ acting repectively on the sections of $S(\mathcal{F})$ and $\bar{S}(\mathcal{F})$, are related, for any vector field $X$ and any spinor field $\Psi$ by

$$
\begin{equation*}
\bar{\nabla}_{X} \bar{\Psi}=\overline{\nabla_{X} \Psi}-\frac{1}{2} \overline{\pi(X) \cdot \operatorname{grad}_{\nabla}(u) \cdot \Psi}-\frac{1}{2} g_{Q}\left(\operatorname{grad}_{\nabla}(u), \pi(X)\right) \bar{\Psi} \tag{3.19}
\end{equation*}
$$

Proof. Let $\left\{E_{a}\right\}$ be an orthonormal basis of $Q$ and denote by $\omega$ and $\bar{\omega}$, the connection forms corresponding to $g_{Q}$ and $\bar{g}_{Q}$. That is, for any vector field $X \in T M$,

$$
\begin{equation*}
\nabla_{X} E_{b}=\sum_{c} \omega_{b c}(\pi(X)) E_{c}, \quad \bar{\nabla}_{X} \bar{E}_{b}=\sum_{c} \bar{\omega}_{b c}(\pi(X)) \bar{E}_{c} \tag{3.20}
\end{equation*}
$$

From (3.18), we have

$$
\begin{equation*}
\bar{\omega}_{b c}(\pi(X))=\omega_{b c}(\pi(X))+g_{Q}\left(\pi(X), E_{c}\right) E_{b}(u)-g_{Q}\left(\pi(X), E_{b}\right) E_{c}(u) . \tag{3.21}
\end{equation*}
$$

Let $\left\{\Psi_{A}\right\}\left(A=1, \cdots, 2^{\left[\frac{q}{2}\right]}\right)$ be a local frame field of $S(\mathcal{F})$. Then the spinor covariant derivative of $\Psi_{A}$ is given ([12]) by

$$
\begin{equation*}
\nabla_{X} \Psi_{A}=\frac{1}{2} \sum_{b<c} \omega_{b c}(\pi(X)) E_{b} \cdot E_{c} \cdot \Psi_{A} \tag{3.22}
\end{equation*}
$$

With respect to $\bar{g}_{Q}$, we have

$$
\begin{aligned}
& \bar{\nabla}_{X} \bar{\Psi}_{A}=\frac{1}{2} \sum_{b<c} \bar{\omega}_{b c}(\pi(X)) \bar{E}_{b} \cdot \bar{E}_{c} \cdot \bar{\Psi}_{A} \\
& =\frac{1}{2} \sum_{b<c}\left\{\omega_{b c}(\pi(X))+g_{Q}\left(\pi(X), E_{c}\right) E_{b}(u)\right. \\
& \left.-g_{Q}\left(\pi(X), E_{b}\right) E_{c}(u)\right\} \bar{E}_{b}-\bar{E}_{c} \cdot \bar{\Psi}_{A} \\
& =\overline{\nabla_{X} \Psi_{A}}-\frac{1}{2} \sum_{b \neq c} g_{Q}\left(\pi(X), E_{c}\right) E_{b}(u) \bar{E}_{c} \cdot \bar{E}_{b} \cdot \bar{\Psi}_{A} \\
& =\overline{\nabla_{X} \Psi_{A}}-\frac{1}{2} \overline{\pi(X) \cdot \operatorname{grad}_{\nabla}(u) \cdot \Psi_{A}}-\frac{1}{2} g_{Q}\left(\operatorname{grad}_{\nabla}(u), \pi(X)\right) \bar{\Psi}_{A} .
\end{aligned}
$$

Let $\bar{D}_{t r}$ be the transversal Dirac operator associated with the metric $\bar{g}_{Q}=e^{2 u} g_{Q}$ and acting on the sections of the foliated spinor bundle $\bar{S}(\mathcal{F})$. Let $\left\{E_{a}\right\}$ be a local frame of $P_{S O}(\mathcal{F})$ and $\left\{\bar{E}_{a}\right\}$ a local frame of $\bar{P}_{S O}(\mathcal{F})$. Locally, $\bar{D}_{t r}$ is expressed by

$$
\begin{equation*}
\bar{D}_{t r} \bar{\Psi}=\sum_{a} \bar{E}_{a} \overline{\nabla_{\bar{E}_{a}}} \bar{\Psi}-\frac{1}{2} \kappa_{\bar{g}} \bar{\Psi}, \tag{3.23}
\end{equation*}
$$

where $\kappa_{\bar{g}}$ is the mean curvature form associated with $\bar{g}_{Q}$, which satisfies $\kappa_{\bar{g}}=e^{-2 u} \kappa$. Using (3.19), we have that for any $\Psi$,

$$
\begin{equation*}
\bar{D}_{t r} \bar{\Psi}=e^{-u}\left\{\overline{D_{t r} \Psi}+\frac{q-1}{2} \overline{\operatorname{grad}_{\nabla}(u) \cdot \Psi}\right\} . \tag{3.24}
\end{equation*}
$$

Now, for any function $f$, we have $D_{t r}(f \Psi)=\operatorname{grad}_{\nabla}(f) \cdot \Psi+f D_{t r} \Psi$. Hence we have

$$
\begin{equation*}
\bar{D}_{t r}(f \bar{\Psi})=e^{-u} \overline{\operatorname{grad}_{\nabla}(f) \cdot \Psi}+f \bar{D}_{t r} \bar{\Psi} . \tag{3.25}
\end{equation*}
$$

From (3.24) and (3.25), we have the following proposition.

Proposition 3.12 Let $\mathcal{F}$ be the transverse spin foliation of codimension $q$. Then the transverse Dirac operators $D_{t r}$ and $\bar{D}_{t r}$ satisfy

$$
\begin{equation*}
\bar{D}_{t r}\left(e^{-\frac{q-1}{2} u} \bar{\Psi}\right)=e^{-\frac{q+1}{2} u} \overline{D_{t r} \Psi} \tag{3.26}
\end{equation*}
$$

for any spinor field $\Psi \in S(\mathcal{F})$.

From Proposition 3.12, if $D_{t r} \Psi=0$, then $\bar{D}_{t r} \bar{\Phi}=0$, where $\Phi=e^{-\frac{q-1}{2} u} \Psi$, and conversely. So we have the following corollary.

Corollary 3.13 On the transverse spin foliation $\mathcal{F}$, the dimension of the space of the foliated harmonic spinors is a transversally conformal invariant.

Let the mean curvature form $\kappa$ of $\mathcal{F}$ be basic- harmonic, i.e., $\kappa \in \Omega_{B}^{1}(\mathcal{F})$ and $\delta_{B} \kappa=0$. Then by direct calculation, we have the Lichnerowicz type formula.

Theorem 3.14 On the transverse spin foliation with the basic harmonic mean curvature form $\kappa$, we have on $\bar{S}(\mathcal{F})$

$$
\begin{equation*}
\bar{D}_{t r}^{2} \bar{\Psi}=\bar{\nabla}_{t r}^{*} \bar{\nabla}_{t r} \bar{\Psi}+\mathcal{R}^{\bar{\nabla}}(\bar{\Psi})+K^{\bar{\nabla}} \bar{\Psi} \tag{3.27}
\end{equation*}
$$

where

$$
\begin{align*}
& \bar{\nabla}_{t r}^{*} \bar{\nabla}_{t r} \bar{\Psi}=-\sum_{a} \bar{\nabla}_{\bar{E}_{a}} \bar{\nabla}_{\bar{E}_{a}} \bar{\Psi}+\bar{\nabla}_{\sum \bar{\nabla}_{\bar{E}_{a}} \bar{E}_{a}} \bar{\Psi}+\bar{\nabla}_{\kappa_{\bar{g}}} \bar{\Psi},  \tag{3.28}\\
& K^{\bar{\nabla}}=\frac{1}{2}(q-2) \kappa_{\bar{g}}(u)+\frac{1}{4}|\bar{\kappa}|^{2},  \tag{3.29}\\
& \mathcal{R}^{\bar{\nabla}}(\bar{\Psi})=\sum_{a<b} \bar{E}_{a} \cdot \bar{E}_{b} \cdot \bar{R}^{S}\left(\bar{E}_{a}, \bar{E}_{b}\right) \bar{\Psi} . \tag{3.30}
\end{align*}
$$

Proof. Fix $x \in M$ and choose a local orthonormal basic frame $\left\{E_{a}\right\}$ satisfying $\left(\nabla E_{a}\right)_{x}=0$ at $x \in M$. Then by definition,

$$
\begin{aligned}
\bar{D}_{t r}^{2} \bar{\Psi}= & \bar{D}_{t r}\left\{\sum_{a} \bar{E}_{a} \div \bar{\nabla}_{\bar{E}_{a}} \bar{\Psi}^{2}-\frac{1}{2} \kappa_{\bar{g}} \div \bar{\Psi}\right\} \\
= & -\sum_{a} \bar{\nabla}_{\bar{E}_{a}} \bar{\nabla}_{\bar{E}_{a}} \bar{\Psi}+\sum_{a<b} \bar{E}_{a} \div \bar{E}_{b} \div \bar{R}^{S}\left(\bar{E}_{a}, \bar{E}_{b}\right) \bar{\Psi} \\
& +\sum_{a<b} \bar{E}_{a} \div \bar{E}_{b} \div \bar{\nabla}_{\left[\bar{E}_{a} \bar{E}_{b} \bar{U}^{\prime}\right.} \bar{\Psi} \pm \sum_{a, b} \bar{E}_{b} \div \bar{\nabla}_{\bar{E}_{b}} \bar{E}_{a} \div \bar{\nabla}_{\bar{E}_{a}} \bar{\Psi} \\
& -\frac{1}{2} \sum_{b} \bar{E}_{b} \div\left(\bar{\nabla}_{\bar{E}_{b}} \kappa_{\bar{g}}\right) \div \bar{\Psi}+\bar{\nabla}_{\kappa_{\bar{g}}} \bar{\Psi}+\frac{1}{4} \kappa_{\bar{g}} \div \kappa_{\bar{g}} \div \bar{\Psi} .
\end{aligned}
$$

From Proposition 3.10, we have

$$
\begin{equation*}
\bar{\nabla}_{\bar{E}_{a}} \bar{E}_{b}=e^{-2 u}\left\{E_{b}(u) E_{a}-\delta_{a b} g r a d_{\nabla}(u)\right\} \tag{3.31}
\end{equation*}
$$

Hence we have

$$
\begin{aligned}
& \sum_{a<b} \bar{E}_{a} \div \bar{E}_{b} \div \bar{\nabla}_{\left[\bar{E}_{a}, \bar{E}_{b}\right]} \bar{\Psi}=e^{-u}\left\{\sum_{a} \overline{E_{a} \cdot \operatorname{grad}_{\nabla}(u)} \div \bar{\nabla}_{\bar{E}_{a}} \bar{\Psi}+\bar{\nabla} \overline{\operatorname{grad}_{\nabla}(u)}\right. \\
& \bar{\Psi}\}, \\
& \sum_{a, b} \bar{E}_{b} \div \bar{\nabla}_{\bar{E}_{b}} \bar{E}_{a} \div \bar{\nabla}_{\bar{E}_{a}} \bar{\Psi}=-e^{-u}\left\{q \bar{\nabla} \overline{\operatorname{grad}_{\nabla}(u)} \bar{\Psi}+\sum_{a} \bar{E}_{a} \div \overline{\operatorname{grad}_{\nabla}(u)} \div \bar{\nabla}_{\bar{E}_{a}} \bar{\Psi}\right\}, \\
& \sum_{a} \bar{E}_{a} \div\left(\bar{\nabla}_{\bar{E}_{a}} \kappa_{\bar{g}}\right) \div \bar{\Psi}=e^{-2 u}\left\{\sum_{a} \overline{E_{a} \cdot \nabla_{E_{a}} \kappa \cdot \Psi}+(2-q) \kappa(u) \bar{\Psi}\right\} .
\end{aligned}
$$

From the above equations, we have

$$
\begin{aligned}
\bar{D}_{t r}^{2} \bar{\Psi}= & -\sum_{a} \bar{\nabla}_{\bar{E}_{a}} \bar{\nabla}_{\bar{E}_{a}} \bar{\Psi}+\bar{\nabla}_{\sum_{a} \bar{\nabla}_{\bar{E}_{a}} \bar{E}_{a}} \bar{\Psi}+\bar{\nabla}_{\kappa_{\bar{g}}} \bar{\Psi} \\
& +\sum_{a<b} \bar{E}_{a} \cdot \bar{E}_{b} \cdot \bar{R}^{S}\left(\bar{E}_{a}, \bar{E}_{b}\right) \bar{\Psi}+\frac{1}{2}(q-2) \kappa_{\bar{g}}(u) \Psi+\frac{1}{4}|\bar{\kappa}|^{2} \bar{\Psi}
\end{aligned}
$$

This completes the proof.
Lemma 3.15 ([12]) Let $\left(M, g_{M}, \mathcal{F}\right)$ be a compact Riemannian manifold with a foliation $\mathcal{F}$ and a bundle-like metric $g_{M}$ with respect to $\mathcal{F}$. Then

$$
\begin{equation*}
\ll \bar{\nabla}_{t r}^{*} \bar{\nabla}_{t r} \bar{\Psi}, \bar{\Phi} \gg_{\bar{g}_{Q}}=\ll \bar{\nabla}_{t r} \bar{\Psi}, \bar{\nabla}_{t r} \bar{\Phi} \gg \bar{g}_{Q} \tag{3.32}
\end{equation*}
$$

for all $\Phi, \Psi \in S(\mathcal{F})$, where $<\bar{\nabla}_{t r} \bar{\Psi}, \bar{\nabla}_{t r} \bar{\Phi}>_{\bar{g}_{Q}}=\sum_{a}<\bar{\nabla}_{\bar{E}_{a}} \bar{\Psi}, \bar{\nabla}_{\bar{E}_{a}} \bar{\Phi}>_{\bar{g}_{Q}}$.
Proof. Fix $x \in M$ and choose an orthonormal basic frame $\left\{E_{a}\right\}$ such that $\left(\nabla E_{a}\right)_{x}=0$ for all $a$. Then we have that at $x$

$$
\begin{equation*}
\bar{\nabla}_{\bar{E}_{a}} \bar{E}_{b}=e^{-2 u}\left\{E_{b}(u) E_{a}-\delta_{a b} \operatorname{grad}_{\nabla}(u)\right\} . \tag{3.33}
\end{equation*}
$$

Hence we have

$$
\begin{aligned}
<\bar{\nabla}_{t r}^{*} \bar{\nabla}_{t r} \bar{\Psi}, \bar{\Phi}>_{\bar{g}_{Q}}= & -\sum_{a}<\bar{\nabla}_{\bar{E}_{a}} \bar{\nabla}_{\bar{E}_{a}} \bar{\Psi}, \bar{\Phi}>_{\bar{g}_{Q}}+<\bar{\nabla}_{\kappa_{\bar{g}}} \bar{\Psi}, \bar{\Phi}>_{\bar{g}_{Q}} \\
& +(1-q) e^{-2 u}<\bar{\nabla}_{g_{r a d_{\nabla}(u)}} \bar{\Psi}, \bar{\Phi}>_{\bar{g}_{Q}} \\
= & -\sum_{a} \bar{E}_{a}<\bar{\nabla}_{\bar{E}_{a}} \bar{\Psi}, \bar{\Phi}>_{\bar{g}_{Q}}+<\bar{\nabla}_{\kappa_{\bar{g}}} \bar{\Psi}, \bar{\Phi}>_{\bar{g}_{Q}} \\
& +(1-q) e^{-2 u}<\bar{\nabla}_{g_{r a d_{\nabla}(u)}} \bar{\Psi}, \bar{\Phi}>_{\bar{g}_{Q}} \\
& +\sum_{a}<\bar{\nabla}_{\bar{E}_{a}} \bar{\Psi}, \bar{\nabla}_{\bar{E}_{a}} \bar{\Phi}>_{\bar{g}_{Q}} \\
= & -\operatorname{div}_{\bar{\nabla}}(V)+\sum_{a}<\bar{\nabla}_{\bar{E}_{a}} \bar{\Psi}, \bar{\nabla}_{\bar{E}_{a}} \bar{\Phi}_{\bar{g}_{Q}} \\
& +<\bar{\nabla}_{\kappa_{\bar{g}}} \bar{\Psi}, \bar{\Phi}>_{\bar{g}_{Q}},
\end{aligned}
$$

where $V \in \Gamma Q \otimes \mathbb{C}$ are defined by $\bar{g}_{Q}(V, Z)=<\bar{\nabla}_{Z} \bar{\Psi}, \bar{\Phi}>_{\bar{g}_{Q}}$ for all $Z \in \Gamma Q$. The last line is proved as follows: At $x \in M$,

$$
\begin{aligned}
\operatorname{div}_{\bar{\nabla}}(V) & =\sum_{a} \bar{g}_{Q}\left(\bar{\nabla}_{\bar{E}_{a}} V, \bar{E}_{a}\right)=\sum_{a} \bar{E}_{a} \bar{g}_{Q}\left(V, \bar{E}_{a}\right)-\bar{g}_{Q}\left(V, \sum_{a} \bar{\nabla}_{\bar{E}_{a}} \bar{E}_{a}\right) \\
& =\sum_{a} \bar{E}_{a}<\bar{\nabla}_{\bar{E}_{a}} \bar{\Psi}, \bar{\Phi}>_{\bar{g}_{Q}}-(1-q) e^{-2 u}<\bar{\nabla}_{\operatorname{grad}_{\nabla}(u)} \bar{\Psi}, \bar{\Phi}>_{\bar{g}_{Q}} .
\end{aligned}
$$

By the transversal divergence theorem on the foliated Riemannian manifold([19,23])

$$
\int_{M} \operatorname{di} v_{\bar{\nabla}}(V) v_{\bar{g}}=\int_{M} \bar{g}_{Q}\left(\kappa_{\bar{g}}, V\right) v_{\bar{g}}=\int_{M}<\bar{\nabla}_{\kappa_{\bar{g}}} \bar{\Psi}, \bar{\Phi}>_{\bar{g}_{Q}} v_{\bar{g}},
$$

where $v_{\bar{g}}$ is the volume form associated to the metric $\bar{g}_{M}=g_{L}+\bar{g}_{Q}$. By integrating, we obtain our result.

## 제주대학교 중앙도서관

## 4 Eigenvalue estimate of the basic Dirac operator

### 4.1 Eigenvalue estimate I

Let $\left(M, \tilde{g}_{M}, \mathcal{F}, S(\mathcal{F})\right)$ be a Riemannian manifold with a transverse spin foliation $\mathcal{F}$ of codimension $q \geq 2$. Let $\tilde{g}_{M}$ be the bundle-like metric for which the mean curvature $\kappa$ is basic-harmonic, i.e., $\kappa \in \Omega_{B}^{1}(\mathcal{F})$ and $\delta_{B} \kappa=0$.

Now, we introduce a new connection $\stackrel{l}{\nabla}$ on $S(\mathcal{F})$ as followings.

Definition 4.1 Let $l$ be a linear symmetric endomorphism of $Q$. For any tangent vector field $X$ and any spinor field $\Psi$, we define the modified connection $\stackrel{l}{\nabla}$ on $S(\mathcal{F})$ by

$$
\begin{equation*}
\stackrel{l}{\nabla}_{X} \Psi=\nabla_{X} \Psi+l(\pi(X)) \cdot \Psi \tag{4.1}
\end{equation*}
$$

where $\pi: T M \rightarrow Q$.
Proposition 4.2 The connection $\stackrel{l}{\nabla}$ is a metric connection. For any tangent vector field $X$, and any spinor fields $\Psi$ and $\Phi$, one has

$$
X g_{Q}(\Psi, \Phi)=g_{Q}\left(\nabla_{X} \Psi, \Phi\right)+g_{Q}\left(\Psi, \stackrel{l}{\nabla}_{X} \Phi\right)
$$

Proof. By the direct calculation, we have

$$
\begin{aligned}
X g_{Q}(\Psi, \Phi) & =g_{Q}\left(\nabla_{X} \Psi, \Phi\right)+g_{Q}\left(\Psi, \nabla_{X} \Phi\right) \\
& =g_{Q}\left(\nabla_{X} \Psi-l(\pi(X)) \cdot \Psi, \Phi\right)+g_{Q}\left(\Psi, \stackrel{l}{\nabla_{X}} \Phi-l(\pi(X)) \cdot \Phi\right) \\
& =g_{Q}\left(\nabla_{X} \Psi, \Phi\right)+g_{Q}\left(\Psi, \stackrel{l}{\nabla_{X}} \Phi\right) .
\end{aligned}
$$

We now define $\nabla^{l}{ }^{*}{ }_{t r} \stackrel{l}{\nabla}{ }_{t r}: \Gamma S(\mathcal{F}) \rightarrow \Gamma S(\mathcal{F})$ as

$$
\begin{equation*}
\stackrel{l}{\nabla}_{t r}^{*} \stackrel{l}{\nabla}_{t r} \Psi=-\sum_{a} \stackrel{l}{\nabla}_{E a, E a}^{2} \Psi+\stackrel{l}{\nabla}_{\kappa} \Psi \tag{4.2}
\end{equation*}
$$

where $\stackrel{l}{\nabla}_{v, w}^{2}=\stackrel{l}{\nabla}{ }_{v} \stackrel{l}{\nabla}_{w}-\stackrel{l}{\nabla} \nabla_{v w}$ for any $v, w \in T M$. Then we have the following lemma.

Lemma 4.3 Let $\left(M, \tilde{g}_{M}, \mathcal{F}\right)$ be a compact Riemanian manifold with a transverse spin foliation $\mathcal{F}$ and a bundle-like metric $\tilde{g}_{M}$ with respect to $\mathcal{F}$. Then

$$
\begin{equation*}
\ll \nabla^{l}{ }_{t r} \stackrel{l}{\nabla}_{t r} \Phi, \Psi \gg=\ll \stackrel{l}{\nabla}_{t r} \Phi, \stackrel{l}{\nabla}_{t r} \Psi \gg \tag{4.3}
\end{equation*}
$$

for all $\Phi, \Psi \in \Gamma S(\mathcal{F})$, where $\ll \Phi, \Psi \gg=\int_{M}<\Phi, \Psi>$ is the (Complex Hermitian) inner product on $S(\mathcal{F})$ 학교 중앙도서관

Proof. Fix $x \in M$ and choose an orthonormal basic frame $\left\{E_{a}\right\}$ such that $\left(\nabla E_{a}\right)_{x}=0$ for all $a$. Then we have at the point $x$ that for any $\Phi, \Psi$,

$$
\begin{aligned}
& <\stackrel{l}{\nabla}^{*}{ }_{t r} \stackrel{l}{\nabla}_{t r} \Phi, \Psi>=-\sum_{a}<\stackrel{l}{\nabla}_{E_{a}} \stackrel{l}{\nabla}_{E_{a}} \Phi, \Psi>+<\stackrel{l}{\nabla}_{\kappa} \Phi, \Psi> \\
& =-\sum_{a} E_{a}<\stackrel{l}{\nabla}_{E_{a}} \Phi, \Psi>+\sum_{a}<\stackrel{l}{\nabla}_{E_{a}} \Phi, \stackrel{l}{\nabla_{E_{a}}} \Psi> \\
& +<\nabla_{\kappa} \Phi, \Psi>+<l(\kappa) \cdot \Phi, \Psi> \\
& =-\sum_{a} E_{a}<\nabla_{E_{a}} \Phi, \Psi>-\sum_{a} E_{a}<l\left(E_{a}\right) \cdot \Phi, \Psi> \\
& +<\nabla_{\kappa} \Phi, \Psi>+<l(\kappa) \cdot \Phi, \Psi> \\
& +\sum_{a}<\stackrel{l}{\nabla}_{E_{a}} \Phi, \stackrel{l}{\nabla_{E_{a}}} \Psi>.
\end{aligned}
$$

So,

$$
\begin{aligned}
<\nabla^{l}{ }_{t r} \stackrel{l}{\nabla}_{t r} \Phi, \Psi>= & -\operatorname{div}_{\nabla}(V)-\operatorname{div}_{\nabla}(W)+\sum_{a}<\nabla_{\nabla_{a}} \Phi, \nabla_{E_{a}} \Psi> \\
& +<\nabla_{\kappa} \Phi, \Psi>+<l(\kappa) \cdot \Phi, \Psi>
\end{aligned}
$$

where $V, W \in \Gamma Q \otimes \mathbb{C}$ are defined by the condition that $g_{Q}(V, Z)=<$ $\nabla_{Z} \Phi, \Psi>$ and $g_{Q}(W, Z)=<l(Z) \cdot \Phi, \Psi>$ for all $Z \in \Gamma Q$. The last line is proved as followings: At $x \in M$,

$$
\begin{align*}
\operatorname{div_{\nabla }(V)} & =\sum_{a} g_{Q}\left(\nabla_{E_{a}} V, E_{a}\right)=\sum_{a} E_{a} g_{Q}\left(V, E_{a}\right) \\
& =\sum_{a} E_{a}<\nabla_{E_{a}} \Phi, \Psi> \tag{4.4}
\end{align*}
$$

Similarly, we have

$$
\begin{align*}
\operatorname{div}_{\nabla}(W) & =\sum_{a} g_{Q}\left(\nabla_{E_{a}} W, E_{a}\right)=\sum_{a} E_{a} g_{Q}\left(W, E_{a}\right)  \tag{4.5}\\
& =\sum_{a} E_{a}<l\left(E_{a}\right) \cdot \Phi, \Psi>.
\end{align*}
$$

By the transversal divergence theorem $([19,23])$ on the foliated Riemannian manifold, we have

$$
\int_{M} \operatorname{div}_{\nabla}(V)=\ll \kappa, V \gg=\ll \nabla_{\kappa} \Phi, \Psi \gg
$$

and

$$
\int_{M} \operatorname{div}_{\nabla}(W)=\ll l(\kappa) \cdot \Phi, \Psi \gg
$$

By integrating,

$$
\begin{aligned}
\int_{M}<\nabla^{l}{ }_{t r} \stackrel{l}{\nabla}_{t r} \Phi, \Psi>= & -\int_{M} \operatorname{div}_{\nabla}(V)-\int_{M} \operatorname{div}_{\nabla}(W) \\
& +\int_{M}<\nabla_{\kappa} \Phi, \Psi>+\int_{M}<l(\kappa) \cdot \Phi, \Psi> \\
& +\int_{M} \sum_{a}<\nabla_{E_{a}}^{l} \Phi, \nabla_{E_{a}}^{l} \Psi>
\end{aligned}
$$

Hence, we obtain our result.
Proposition 4.4 For any linear symmetric endomorphism $l$ of $Q$, and for any spinor field $\Psi$, the following identity holds:

$$
\begin{align*}
\left|\stackrel{l}{\nabla}_{t r} \Psi\right|^{2} & :=\sum_{a}<\stackrel{l}{\nabla}_{E_{a}} \Psi, \stackrel{l}{\nabla_{E_{a}}} \Psi>  \tag{4.6}\\
& =\left|\nabla_{t r} \Psi\right|^{2}-2 R e \sum_{a}<l\left(E_{a}\right) \cdot \nabla_{E_{a}} \Psi, \Psi>+|l|^{2}|\Psi|^{2} .
\end{align*}
$$

Proof. Fix $x \in M$ and choose an orthonormal basic frame $\left\{E_{a}\right\}$ such that $\left(\nabla E_{a}\right)_{x}=0$ for all $a$. Then we have at the point $x$ that for any $\Psi$,

$$
\begin{aligned}
\left|\stackrel{l}{\nabla}_{t r} \Psi\right|^{2}:= & \sum_{a}<\nabla_{E_{a}} \Psi, \stackrel{l}{\nabla_{E_{a}}} \Psi> \\
= & \sum_{a}<\nabla_{E_{a}} \Psi+l\left(\pi\left(E_{a}\right)\right) \cdot \Psi, \nabla_{E_{a}} \Psi+l\left(\pi\left(E_{a}\right)\right) \cdot \Psi> \\
= & \sum_{a}<\nabla_{E_{a}} \Psi, \nabla_{E_{a}} \Psi>+\sum_{a}<\nabla_{E_{a}} \Psi, l\left(E_{a}\right) \cdot \Psi> \\
& +\sum_{a}<l\left(E_{a}\right) \cdot \Psi, \nabla_{E_{a}} \Psi>+\sum_{a}<l\left(E_{a}\right) \cdot \Psi, l\left(E_{a}\right) \cdot \Psi> \\
= & \left|\nabla_{t r} \Psi\right|^{2}-\sum_{a}\left\{<l\left(E_{a}\right) \nabla_{E_{a}} \Psi, \Psi>+<\Psi, l\left(E_{a}\right) \nabla_{E_{a}} \Psi>\right\} \\
& +|l|^{2}|\Psi|^{2} \\
= & \left|\nabla_{t r} \Psi\right|^{2}-2 R e \sum_{a}<l\left(E_{a}\right) \cdot \nabla_{E_{a}} \Psi, \Psi>+|l|^{2}|\Psi|^{2} .
\end{aligned}
$$

We now show that for an appropriate choice of the symmetric endomorphism $l$, one gets a sharp estimate of the first eigenvalue of the basic Dirac operator on compact foliated Riemannian manifolds. For this, we need the following:

Definition 4.5 On the complement of the set of zeroes of a spinor field $\Psi \in \Gamma_{B} S(\mathcal{F})$, define for any tangent vector fields $X$ and $Y$, the symmetric
bilinear tensor $F_{\Psi}$ by

$$
\begin{equation*}
F_{\Psi}(X, Y)=\frac{1}{2} R e<\pi(X) \cdot \nabla_{Y} \Psi+\pi(Y) \cdot \nabla_{X} \Psi, \Psi /|\Psi|^{2}> \tag{4.7}
\end{equation*}
$$

where $\pi: T M \rightarrow Q$.
Let $l_{\Psi}$ be a symmetric linear map associated to $F_{\Psi}$. Namely, it follows that for any $X, Y \in \Gamma Q$

$$
\begin{equation*}
F_{\Psi}(X, Y)=g_{Q}\left(l_{\Psi}(X), Y\right) . \tag{4.8}
\end{equation*}
$$

Since $<\kappa \cdot \Psi, \Psi>$ is pure imaginary, we have

$$
\begin{aligned}
\operatorname{tr} l_{\Psi} & =\sum_{a} R e<E_{a} \cdot \nabla_{E_{a}} \Psi, \Psi /|\Psi|^{2}> \\
& =R e<D_{b} \Psi, \Psi /|\Psi|^{2}>
\end{aligned}
$$

It is trivial that if $\Psi$ satisfies $D_{b} \Psi=\lambda \Psi$, then

$$
\begin{align*}
& \text { 제주대하ㄱㅛㅛ 중앙도서관 }  \tag{4.9}\\
& \operatorname{tr} l_{\Psi}=\lambda .
\end{align*}
$$

On the other hand, since $l$ is linear symmetric endomorphism, we have

$$
\begin{align*}
\operatorname{Re} \sum_{a}<l\left(E_{a}\right) \cdot \nabla_{E_{a}} \Psi, \Psi> & =\sum_{a}<l\left(E_{a}\right), l_{\Psi}\left(E_{a}\right)>|\Psi|^{2}  \tag{4.10}\\
& =<l, l_{\Psi}>|\Psi|^{2} .
\end{align*}
$$

From (4.6) and (4.10), we have the following equation;

$$
\begin{equation*}
\left|\stackrel{l}{\nabla}_{t r} \Psi\right|^{2}=\left|\nabla_{t r} \Psi\right|^{2}-f(l)|\Psi|^{2} \tag{4.11}
\end{equation*}
$$

where $f(l)=2<l, l_{\Psi}>-|l|^{2}$. Note that $f(l)$ has maximum value $\left|l_{\Psi}\right|^{2}$ at $l=l_{\Psi}$ because $l$ is linear endomorphism. From (3.14) and (4.11), we have that for any eigenspinor $\Psi$ corresponding to an eigenvalue $\lambda$

$$
\begin{equation*}
\int_{M}\left|\stackrel{l}{\Psi}_{\nabla_{t r}} \Psi\right|^{2}=\int_{M}\left\{\lambda^{2}-\left(\frac{1}{4} K^{\sigma}+\left|l_{\Psi}\right|^{2}\right)\right\}|\Psi|^{2} \tag{4.12}
\end{equation*}
$$

Hence we have the following theorem.

Theorem 4.6 $\operatorname{Let}\left(M, \tilde{g}_{M}, \mathcal{F}\right)$ be a compact Riemannian manifold with a transverse spin foliation $\mathcal{F}$ and a bundle-like metric $\tilde{g}_{M}$. Then any eigenvalue $\lambda$ of the basic Dirac operator $D_{b}$ corresponding to the eigenspinor $\Psi \in \Gamma S(\mathcal{F})$ satisfies

$$
\begin{equation*}
\lambda^{2} \geq \inf _{M}\left(\frac{1}{4} K^{\sigma}+\left|l_{\Psi}\right|^{2}\right) \tag{4.13}
\end{equation*}
$$

where $K^{\sigma}=\sigma^{\nabla}+|\kappa|^{2}$.

Moreover, by the Cauchy-Schwarz inequality we have

$$
\begin{aligned}
\frac{1}{q}\left(\operatorname{tr} l_{\Psi}\right)^{2} & =\frac{1}{q}\left(\sum_{a}<l_{\Psi}\left(E_{a}\right), E_{a}>\right)^{2} \\
& \leq \frac{1}{q} \cdot q \sum_{a}\left|l_{\Psi}\left(E_{a}\right)\right|^{2}\left|E_{a}\right|^{2} \\
& =\sum_{a}\left|l_{\Psi}\left(E_{a}\right)\right|^{2}=\left|l_{\Psi}\right|^{2} . \\
& \text { 제주 }{ }_{a} \text { 아ㅍㅔㅜㅇㅇㅏㅏ더과 }
\end{aligned}
$$

Hence from (4.9), we have

$$
\begin{equation*}
\left|l_{\Psi}\right|^{2} \geq \frac{\lambda^{2}}{q} \tag{4.14}
\end{equation*}
$$

Hence we have the following corollary.

Corollary 4.7 (cf.[11]) Under the same conditions as in Theorem 4.6, one has

$$
\begin{equation*}
\lambda^{2} \geq \frac{q}{4(q-1)} \inf _{M} K^{\sigma} \tag{4.15}
\end{equation*}
$$

### 4.2 Eigenvalue estimate II

Now, we introduce a connection $\stackrel{l}{\nabla}$ on $\bar{S}(\mathcal{F})$, as

$$
\begin{equation*}
\stackrel{l}{\nabla}_{X} \bar{\Psi}=\bar{\nabla}_{X} \bar{\Psi}+l(\pi(X)) \div \bar{\Psi} \quad \text { for } X \in T M \tag{4.16}
\end{equation*}
$$

where $l$ is a linear symmetric endomorphism on $Q$. Trivially, the connection $\stackrel{l}{\nabla}$ is a metric connection and $\stackrel{l}{\nabla}^{*} \stackrel{l}{\nabla}_{t r}$ is positive definite.

Lemma 4.8 On the foliated spinor bundle $\bar{S}(\mathcal{F})$, we have

$$
\ll \bar{\nabla}^{*}{ }_{t r} \stackrel{l}{\nabla}_{t r} \bar{\Psi}, \bar{\Phi} \gg_{\bar{g}_{Q}}=\ll \stackrel{l}{\nabla}_{t r} \bar{\Psi}, \stackrel{l}{\nabla}_{t r} \bar{\Phi} \gg_{\bar{g}_{Q}}
$$

for all $\Psi, \Phi \in \Gamma S(\mathcal{F})$.
On the other hand, we obtain the following lemma.
Lemma 4.9 On the spinor bundle $\bar{S}(\mathcal{F})$ associated with the metric $\bar{g}_{Q}=$ $e^{2 u} g_{Q}$

$$
\begin{equation*}
F_{\bar{\Phi}}(\bar{X}, \bar{Y})=e^{-u} F_{\Phi}(X, Y)=e^{-u} F_{\Psi}(X, Y) \tag{4.17}
\end{equation*}
$$

for any $X, Y \in \Gamma Q$, where $\Phi=e^{-\frac{q-1}{2} u} \Psi$.
Proof. By definition, we have

$$
\begin{aligned}
F_{\bar{\Phi}}(\bar{X}, \bar{Y})= & \frac{1}{2} R e<\bar{X} \cdot \bar{\nabla}_{\bar{Y}} \bar{\Phi}+\bar{Y}^{-} \bar{\nabla}_{\bar{X}} \bar{\Phi}, \bar{\Phi} /|\bar{\Phi}|_{\bar{g}_{Q}}^{2}>_{\bar{g}_{Q}} \\
= & \frac{1}{2} e^{-u} R e<\bar{X}^{-} \cdot\left\{\overline{\nabla_{Y} \Phi}-\frac{1}{2} \overline{Y \cdot g r a d_{\nabla}(u) \cdot \Phi}\right. \\
& \left.-\frac{1}{2} g_{Q}\left(\operatorname{grad}_{\nabla}(u), Y\right) \bar{\Phi}\right\} \\
& +\bar{Y} \cdot\left\{\overline{\nabla_{X} \Phi}-\frac{1}{2} \overline{X \cdot \operatorname{grad}}{ }_{\nabla}(u) \cdot \Phi\right. \\
& \left.-\frac{1}{2} g_{Q}\left(\operatorname{grad}_{\nabla}(u), X\right) \bar{\Phi}\right\}, \bar{\Phi} /|\bar{\Phi}|_{\bar{g}_{Q}}^{2}>_{\bar{g}_{Q}} .
\end{aligned}
$$

From (3.16) and (3.17), we have

$$
\begin{aligned}
F_{\bar{\Phi}}(\bar{X}, \bar{Y})= & \frac{1}{2} e^{-u} R e<\overline{X \cdot \nabla_{Y} \Phi+Y \cdot \nabla_{X} \Phi} \\
& \frac{g_{Q}(X, Y) \operatorname{grad}_{\nabla}(u) \cdot \Phi}{}-\frac{1}{2} g_{Q}\left(\operatorname{grad}_{\nabla}(u), Y\right) \overline{X \cdot \Phi} \\
& -\frac{1}{2} g_{Q}\left(\operatorname{grad}_{\nabla}(u), X\right) \overline{Y \cdot \Phi}, \bar{\Phi} /|\bar{\Phi}|_{\bar{g}_{Q}}^{2}>_{\bar{g}_{Q}} .
\end{aligned}
$$

Since $\tilde{I}_{u}$ is isometry and $<X \cdot \Psi, \Psi>$ is pure imaginary, it follows that

$$
\begin{aligned}
F_{\bar{\Phi}}(\bar{X}, \bar{Y}) & =\frac{1}{2} e^{-u} R e<X \cdot \nabla_{Y} \Phi+Y \cdot \nabla_{X} \Phi, \Phi /|\Phi|^{2}> \\
& =e^{-u} F_{\Phi}(X, Y)
\end{aligned}
$$

On the other hand, since $\Phi=e^{-(q-1) u / 2} \Psi$ by direct calculation we have

$$
F_{\Phi}(X, Y)=F_{\Psi}(X, Y)
$$

Hence the proof is completed.
From (4.17), we have the following identity

$$
\begin{equation*}
l_{\bar{\Phi}}=e^{-u} l_{\Phi}=e^{-u} l_{\Psi} . \tag{4.18}
\end{equation*}
$$

From (4.18), we have the following proposition (see [10] for the details).
Proposition 4.10 The following relations hold:
where $\Phi=e^{-\frac{q-1}{2} u} \Psi$.
By using the connection $\stackrel{l_{\underline{\Phi}}}{\nabla}$, we can obtain the following equation

$$
\begin{aligned}
\left|l_{\bar{\nabla}_{t r}} \bar{\Phi}\right|_{\bar{g}_{Q}}^{2} & =\sum_{a}<l_{\overline{\bar{\nabla}_{\bar{E}_{a}}}} \bar{\Phi},{\stackrel{l \overline{\bar{\Phi}}}{\bar{E}_{a}}}^{\bar{\Phi}^{\prime}}>_{\bar{g}_{Q}} \\
& =\left|\bar{\nabla}_{t r} \bar{\Phi}\right|^{2}-2 R e \sum_{a}<l_{\bar{\Phi}}\left(\bar{E}_{a}\right) \cdot \bar{\nabla}_{\bar{E}_{a}} \bar{\Phi}, \bar{\Phi}>_{\bar{g}_{Q}}+\left.\left|l_{\bar{\Phi}}\right|^{2}\right|_{\bar{g}_{Q}}|\bar{\Phi}|_{\bar{g}_{Q}}
\end{aligned}
$$

Since

$$
\begin{aligned}
& R e \sum_{a}<l_{\bar{\Phi}}\left(\bar{E}_{a}\right) \cdot \bar{\nabla}_{\bar{E}_{a}} \bar{\Phi}, \bar{\Phi}>_{\bar{g}_{Q}} \\
& \quad=\frac{1}{2} R e \sum_{a}<\bar{E}_{a} \div \bar{\nabla}_{l_{\bar{\Phi}}\left(\bar{E}_{a}\right)} \bar{\Phi}+l_{\bar{\Phi}}\left(\bar{E}_{a}\right) \div \bar{\nabla}_{\bar{E}_{a}} \bar{\Phi}, \bar{\Phi}>_{\bar{g}_{Q}} \\
& \quad=\sum_{a} F_{\bar{\Phi}}\left(\bar{E}_{a}, l_{\bar{\Phi}}\left(\bar{E}_{a}\right)\right)|\bar{\Phi}|_{\bar{g}_{Q}}^{2}=\sum_{a}<l_{\bar{\Phi}}\left(\bar{E}_{a}\right), l_{\bar{\Phi}}\left(\bar{E}_{a}\right)>_{\bar{g}_{Q}}|\bar{\Phi}|_{\bar{g}_{Q}}^{2} \\
& \quad=\left|l_{\bar{\Phi}}\right| \overline{\bar{g}}_{Q} \\
& 2
\end{aligned} \bar{\Phi}_{\bar{g}_{Q}}^{2} .
$$

Summing up the above equations, we have

$$
\begin{equation*}
\left.\left|l_{\overline{\bar{\Phi}}}^{t r} \bar{\Phi}_{\overline{\bar{g}}_{Q}}^{2}=\left|\bar{\nabla}_{t r} \bar{\Phi}\right|^{2}-\left|l_{\bar{\Phi}}\right|_{\bar{g}_{Q}}^{2}\right| \bar{\Phi}\right|_{\bar{g}_{Q}} ^{2} \tag{4.20}
\end{equation*}
$$

Hence we have from (3.27), (3.32) and (4.20)

$$
\begin{equation*}
\left.\int_{M}| |_{t r}^{\bar{\nabla}_{t r}} \bar{\Phi}\right|_{\bar{g}_{Q}} ^{2}=\int_{M}\left\{<\bar{D}_{t r}^{2} \bar{\Phi}, \bar{\Phi}>_{\bar{g}_{Q}}-\frac{1}{4}<K_{\sigma}^{\bar{\nabla}} \bar{\Phi}, \bar{\Phi}>_{\bar{g}_{Q}}-\left|l_{\bar{\Phi}}\right|_{\bar{g}_{Q}}^{2}|\bar{\Phi}|_{\bar{g}_{Q}}^{2}\right\} . \tag{4.21}
\end{equation*}
$$

where $K_{\sigma}^{\overline{\bar{\nabla}}}=\sigma^{\bar{\nabla}}+4 K^{\bar{\nabla}}$.
Let $D_{b} \Psi=\lambda \Psi(\Psi \neq 0)$. From Proposition 3.12, we have

$$
\begin{equation*}
\bar{D}_{b} \bar{\Phi}=\lambda e^{-u} \bar{\Phi} \tag{4.22}
\end{equation*}
$$

where $\Phi=e^{-\frac{q-1}{2} u} \Psi$. From (4.19) and (4.22), we have that for any eigenspinor $\Psi$ corresponding to the eigenvalue $\lambda$

$$
\begin{align*}
\int \mid{\stackrel{\nabla_{t r}}{t}}_{\left.\bar{\Phi}\right|_{\bar{g}_{Q}} ^{2}} & =\int e^{-2 u}\left\{\lambda^{2}-\frac{1}{4}\left(e^{2 u} K_{\sigma}^{\text {젱 }}+4\left|l_{\Psi}\right|^{2}\right)\right\}|\bar{\Phi}|_{\bar{g}_{Q}}^{2}  \tag{4.23}\\
& \leq \int e^{-2 u}\left\{\lambda^{2}-\frac{1}{4} \inf _{M}\left(e^{2 u} K_{\sigma}^{\bar{\nabla}}+4\left|l_{\Psi}\right|^{2}\right)\right\}|\bar{\Phi}|_{\bar{g}_{Q}}^{2}
\end{align*}
$$

where $\Phi=e^{-\frac{q-1}{2} u} \Psi$. From (4.23), we have that $\lambda^{2} \geq \frac{1}{4} \inf _{M}\left(e^{2 u} K_{\sigma}^{\bar{\nabla}}+\right.$ $\left.4\left|l_{\Psi}\right|^{2}\right)$. Hence we have the following theorem.

Theorem 4.11 Let $\left(M, \tilde{g}_{M}, \mathcal{F}\right)$ be a compact Riemannian manifold with a transverse spin foliation $\mathcal{F}$ of codimension $q \geq 2$ and bundle-like metric $\tilde{g}_{M}$. Assume that $K_{\sigma}^{\bar{\nabla}} \geq 0$ for some transversally conformal metric $\bar{g}_{Q}=e^{2 u} g_{Q}$. Then any eigenvalue $\lambda$ of the basic Dirac operator $D_{b}$ corresponding to the eigenspinor $\Psi \in \Gamma S(\mathcal{F})$ satisfies

$$
\begin{equation*}
\lambda^{2} \geq \frac{1}{4} \sup _{u} \inf _{M}\left(e^{2 u} K_{\sigma}^{\bar{\nabla}}+4\left|l_{\Psi}\right|^{2}\right), \tag{4.24}
\end{equation*}
$$

where $K_{\sigma}^{\bar{\nabla}}=\sigma^{\bar{\nabla}}+|\bar{\kappa}|^{2}+2(q-2) \kappa_{\bar{g}}(u)$.

From (4.14), we have the following corollary.
Corollary 4.12 (cf.[12]) Under the same assumption as in Theorem 4.11, we have

$$
\begin{equation*}
\lambda^{2} \geq \frac{q}{4(q-1)} \sup _{u} \inf _{M}\left(e^{2 u} K_{\sigma}^{\bar{\sigma}}\right) . \tag{4.25}
\end{equation*}
$$

The transversal Ricci curvature $\rho^{\bar{\nabla}}$ of $\bar{g}_{Q}=e^{2 u} g_{Q}$ and the transversal scalar curvature $\sigma^{\bar{\nabla}}$ of $\bar{g}_{Q}$ are related to the transversal Ricci curvature $\rho^{\nabla}$ of $g_{Q}$ and the transversal scalar curvature $\sigma^{\nabla}$ of $g_{Q}$ by the following lemma.

Lemma 4.13 On a Riemannian foliation $\mathcal{F}$, we have that for any $X \in$ $Q$,

$$
\begin{equation*}
e^{2 u} \sigma^{\bar{\nabla}}=\sigma^{\nabla}+(q-1)(2-q)\left|\operatorname{grad}_{\nabla}(u)\right|^{2}+2(q-1)\left\{\Delta_{B} u-\kappa(u)\right\} . \tag{4.27}
\end{equation*}
$$

Proof. Let $x \in M$ and choose an orthonormal basic frame $\left\{E_{a}\right\}$ with the property that $\left(\nabla E_{a}\right)_{x}=0$ for all $a$. Then

$$
\begin{aligned}
\rho^{\bar{\nabla}}(X) & =\sum_{a} R^{\bar{\nabla}}\left(X, \bar{E}_{a}\right) \bar{E}_{a} \\
& =\sum_{a} \bar{\nabla}_{X} \bar{\nabla}_{E_{a}} \bar{E}_{a}-\sum_{a} \bar{\nabla}_{\bar{E}_{a}} \bar{\nabla}_{X} \bar{E}_{a}-\sum_{a} \bar{\nabla}_{\left[X, \bar{E}_{a}\right]} \bar{E}_{a} .
\end{aligned}
$$

By direct calculation, we have

$$
\begin{aligned}
e^{2 u} \sum_{a} \bar{\nabla}_{X} \bar{\nabla}_{\bar{E}_{a}} \bar{E}_{a}= & (1-q)\left\{\nabla_{X} \operatorname{grad}_{\nabla}(u)+\left|\operatorname{grad}_{\nabla}(u)\right|^{2} X\right. \\
& \left.-2 X(u) \operatorname{grad}_{\nabla}(u)\right\}+\sum_{a} \nabla_{X} \nabla_{E_{a}} E_{a}
\end{aligned}
$$

$$
\begin{align*}
& e^{2 u} \rho^{\bar{\nabla}}(X)=\rho^{\nabla}(X)+(2-q) \nabla_{X} \operatorname{grad}_{\nabla}(u)+(2-q)\left|\operatorname{grad}_{\nabla}(u)\right|^{2} X \\
& +(q-2) X(u) \operatorname{rad} d_{\nabla}(u)+\left\{\Delta_{B} u-\kappa(u)\right\} X . \tag{4.26}
\end{align*}
$$

Similarly,

$$
\begin{aligned}
e^{2 u} \sum_{a} \bar{\nabla}_{\bar{E}_{a}} \bar{\nabla}_{X} \bar{E}_{a}= & \sum_{a} \nabla_{E_{a}} \nabla_{X} E_{a}+\sum_{a} E_{a} E_{a}(u) X \\
& +\nabla_{\operatorname{grad}_{\nabla}(u)} X-\sum_{a} g\left(\nabla_{E_{a}} X, E_{a}\right) \operatorname{grad}_{\nabla}(u) \\
& -\nabla_{X} \operatorname{grad}_{\nabla}(u)-\left|\operatorname{grad}_{\nabla}(u)\right|^{2} X-X(u) \operatorname{grad}_{\nabla}(u)
\end{aligned}
$$

and

$$
\begin{aligned}
e^{2 u} \sum_{a} \bar{\nabla}_{\left[X, \bar{E}_{a}\right]} \bar{E}_{a}= & \sum_{a} \nabla_{\left[X, E_{a}\right]} E_{a}+X(u)(q-1) \operatorname{grad}_{\nabla}(u) \\
& -\nabla_{\operatorname{grad}(u)} X+\sum_{a} g\left(\nabla_{E_{a}} X, E_{a}\right) \operatorname{grad}_{\nabla}(u) .
\end{aligned}
$$

Since $\Delta_{B} u=\delta_{B} d_{B} u=-\sum_{a} E_{a} E_{a}(u)+i(\kappa) d_{B} u$, the above equations give (4.26).

On the other hand,

$$
\sigma^{\bar{\nabla}}=\sum_{a} \bar{g}_{Q}\left(\rho^{\overline{ }}\left(\bar{E}_{a}\right), \bar{E}_{a}\right)=\sum_{a} g_{Q}\left(\rho^{\bar{\nabla}}\left(E_{a}\right), E_{a}\right) .
$$

From (4.26) we have

$$
\begin{aligned}
e^{2 u} \sigma^{\bar{\nabla}}= & \sum_{a} g_{Q}\left(e^{2 u} \rho^{\bar{\nabla}}\left(E_{a}\right), E_{a}\right) \\
= & \sigma^{\nabla}+(2-q) \sum_{a} g_{Q}\left(\nabla_{E_{a}} g r a d_{\nabla}(u), E_{a}\right) \\
& +(q-1)(2-q)\left|\operatorname{grad}_{\nabla}(u)\right|^{2}+q\left\{\Delta_{B} u-\kappa(u)\right\} .
\end{aligned}
$$

Since $\sum_{a} g_{Q}\left(\nabla_{E_{a}} \operatorname{grad}_{\nabla}(u), E_{a}\right)=\sum_{a} E_{a} E_{a}(u)=-\Delta_{B} u+\kappa(u)$, we have

$$
e^{2 u} \sigma^{\bar{\nabla}}=\sigma^{\nabla}+(q-1)(2-q)\left|\operatorname{grad}_{\nabla}(u)\right|^{2}+2(q-1)\left\{\Delta_{B} u-\kappa(u)\right\}
$$

which proves (4.27).

Since $K_{\sigma}^{\bar{\nabla}}=\sigma^{\bar{\nabla}}+|\bar{\kappa}|^{2}+2(q-2) \kappa_{\bar{g}}(u)$, from (4.27), we have

$$
\begin{equation*}
e^{2 u} K_{\sigma}^{\bar{\nabla}}=\sigma^{\nabla}+|\kappa|^{2}+2(q-1) \Delta_{B} u+(q-1)(2-q)\left|\operatorname{grad}_{\nabla}(u)\right|^{2}-2 \kappa(u) . \tag{4.28}
\end{equation*}
$$

On the other hand, for $q \geq 3$, if we choose the positive function $h$ by $u=\frac{2}{q-2} \ln h$, then we have

$$
\begin{align*}
& \Delta_{B} u=\frac{2}{q-2}\left\{h^{-2}\left|\operatorname{grad}_{\nabla}(h)\right|^{2}+h^{-1} \Delta_{B} h\right\},  \tag{4.29}\\
& \left|\operatorname{grad}_{\nabla}(u)\right|^{2}=\left(\frac{2}{q-2}\right)^{2} h^{-2}\left|\operatorname{grad}_{\nabla}(h)\right|^{2} . \tag{4.30}
\end{align*}
$$

From (4.28), (4.29) and (4.30), we have

$$
\begin{equation*}
e^{2 u} K_{\sigma}^{\bar{\nabla}}=h^{\frac{4}{q-2}} K_{\sigma}^{\bar{\nabla}}=h^{-1} Y_{b} h+|\kappa|^{2}-\frac{4}{q-2} h^{-1} \kappa(h), \tag{4.31}
\end{equation*}
$$

where
which is called a basic Yamabe operator of $\mathcal{F}$.
Now we put $\mathcal{K}_{u}=\left\{u \in \Omega_{B}^{0}(\mathcal{F}) \mid \kappa(u)=0\right\}$. If we choose $u \in \mathcal{K}_{u}$, then $\kappa(h)=0=\kappa(u)$. From (4.28) and (4.31), we have
$e^{2 u} K_{\sigma}^{\bar{\nabla}}=K^{\sigma}+2(q-1) \Delta_{B} u+(q-1)(2-q)\left|\operatorname{grad}_{\nabla}(u)\right|^{2}=h^{-1} Y_{b} h+|\kappa|^{2}$,
where $K^{\sigma}=\sigma^{\nabla}+|\kappa|^{2}$. Assume that the transversal scalar curvature $\sigma^{\nabla}$ is non-negative. Then the eigenfunction $h_{1}$ associated to the first eigenvalue $\mu_{1}$ of $Y_{b}$ can be chosen to be positive and then $\mu_{1}$ is non-negative. Thus

$$
\begin{equation*}
h_{1}^{-1} Y_{b} h_{1}=\mu_{1} . \tag{4.34}
\end{equation*}
$$

Since $\sup \inf \left\{h^{-1} Y_{b} h\right\} \geq \mu_{1}$, we have from (4.24) in Theorem 4.11 the following Theorem.

Theorem 4.14 Let $\left(M, \tilde{g}_{M}, \mathcal{F}\right)$ be a compact Riemannian manifold with a transverse spin foliation $\mathcal{F}$ of codimension $q \geq 3$ and bundle-like metric $\tilde{g}_{M}$. Then, any eigenvalue $\lambda$ of the basic Dirac operator corresponding to the eigenspinor $\Psi \in \Gamma S(\mathcal{F})$ satisfies

$$
\begin{equation*}
\lambda^{2} \geq \frac{1}{4}\left(\mu_{1}+\inf _{M}|\kappa|^{2}\right)+\inf _{M}\left|l_{\Psi}\right|^{2} \tag{4.35}
\end{equation*}
$$

where $\mu_{1}$ is the first eigenvalue of the basic Yamabe operator of $\mathcal{F}$

$$
\begin{equation*}
Y_{b}=4 \frac{q-1}{q-2} \Delta_{B}+\sigma^{\nabla} \tag{4.36}
\end{equation*}
$$

acting on functions, and $l_{\Psi}$ is the field of symmetric endomorphism associated with $F_{\Psi}$.

From (4.14), we have the following corollary

Corollary 4.15 (cf.[12]) Let $\left(M, \tilde{g}_{M}, \mathcal{F}\right)$ be a compact Riemannian manifold with a transverse spin foliation $\mathcal{F}$ of codimension $q \geq 3$ and bundlelike metric $\tilde{g}_{M}$. If the transversal scalar curvature satisfies $\sigma^{\nabla} \geq 0$, then any eigenvalue $\lambda$ of $D_{b}$ satisfies

$$
\begin{equation*}
\lambda^{2} \geq \frac{q}{4(q-1)}\left(\mu_{1}+\inf |\kappa|^{2}\right) \tag{4.37}
\end{equation*}
$$

Remark. Since $\mu_{1} \geq \inf \sigma^{\nabla}$, the inequality (4.37) is a sharper estimate than the previous one (1.4). Moreover, Corollary 4.15 is a specialization of the result on an ordinary manifold by O. Hijazi ([10]) to the case of Riemannian foliations.

## 5 The limiting cases

In this chapter, we study the limiting foliations of (4.13) and (4.35).

### 5.1 The limiting case I

For any linear symmetric endomorphism $l$ of $Q$, we define $\operatorname{Ric}_{\nabla}^{l}: \Gamma Q \otimes$ $S(\mathcal{F}) \rightarrow S(\mathcal{F})([10,11])$ by

$$
\begin{equation*}
\operatorname{Ric}_{\nabla}^{l}(X \otimes \Psi)=\sum_{a} E_{a} \cdot R^{l}\left(X, E_{a}\right) \Psi \tag{5.1}
\end{equation*}
$$

where $R^{l}$ is the curvature tensor with respect to $\nabla_{\nabla}^{l}$ defined by (4.1).
Lemma 5.1 On the transverse spin foliation $\mathcal{F}$, we have that for $X \in$ $\Gamma Q$ and $\Psi \in \Gamma S(\mathcal{F})$
$\operatorname{Ric}_{\nabla}^{l}(X \otimes \Psi)=-\frac{1}{2} \rho^{\nabla}(X) \cdot \Psi+\sum_{a}\left\{E_{a} \cdot d l\left(X, E_{a}\right)+E_{a} \cdot\left[l(X), l\left(E_{a}\right)\right]\right\} \cdot \Psi$,
where $[X, Y]=X \cdot Y-Y \cdot X$ and $d l(X, Y)=\left(\nabla_{X} l\right)(Y)-\left(\nabla_{Y} l\right)(X)$.
Proof. Fix $x \in M$ and choose an orthonormal basic frame $\left\{E_{a}\right\}$ such that $\left(\nabla E_{a}\right)_{x}=0$ for all $a$. For $X \in \Gamma Q$, we have

$$
\begin{aligned}
\stackrel{l}{\nabla} \stackrel{l}{\nabla}_{E_{a}} \Psi= & \stackrel{l}{\nabla} X\left\{\nabla_{E_{a}} \Psi+l\left(E_{a}\right) \cdot \Psi\right\} \\
= & \nabla_{X} \nabla_{E_{a}} \Psi+\nabla_{X} l\left(E_{a}\right) \cdot \Psi+l\left(E_{a}\right) \cdot \nabla_{X} \Psi+l(X) \cdot \nabla_{E_{a}} \Psi \\
& +l(X) \cdot l\left(E_{a}\right) \cdot \Psi .
\end{aligned}
$$

Similarly,

$$
\begin{aligned}
\stackrel{l}{\nabla_{E_{a}}} \stackrel{l}{\nabla} \mathrm{X} \Psi= & \nabla_{E_{a}} \nabla_{X} \Psi+\nabla_{E_{a}} l(X) \cdot \Psi+l(X) \cdot \nabla_{E_{a}} \Psi+l\left(E_{a}\right) \cdot \nabla_{X} \Psi \\
& +l\left(E_{a}\right) \cdot l(X) \cdot \Psi
\end{aligned}
$$

From the above equations, we have

$$
\begin{aligned}
\operatorname{Ric}_{\nabla}^{l}(X \otimes \Psi)= & \sum_{a} E_{a} \cdot R^{l}\left(X, E_{a}\right) \Psi \\
= & \sum_{a} E_{a} \cdot\left\{R^{S}\left(X, E_{a}\right) \Psi+\left(\left(\nabla_{X} l\right) E_{a}-\left(\nabla_{E_{a}} l\right) X\right) \cdot \Psi\right. \\
& \left.+\left(l(X) \cdot l\left(E_{a}\right)-l\left(E_{a}\right) \cdot l(X)\right) \cdot \Psi\right\} \\
= & \sum_{a} E_{a} \cdot R^{S}\left(X, E_{a}\right) \Psi+\sum_{a} E_{a} \cdot d l\left(X, E_{a}\right) \cdot \Psi \\
& +\sum_{a} E_{a} \cdot\left[l(X), l\left(E_{a}\right)\right] \cdot \Psi
\end{aligned}
$$

Hence Theorem 3.7 gives our proof.
By the definition of Clifford multiplication, we have

$$
\begin{equation*}
\sum_{a} E_{a} \cdot d l\left(X, E_{a}\right)=\sum_{a} E_{a} \wedge d l\left(X, E_{a}\right)-\left\{X(\operatorname{tr} l)-\left(d i v_{\nabla} l\right)(X)\right\} \tag{5.3}
\end{equation*}
$$

where $\left(d i v_{\nabla} l\right)(X)=\sum_{a} g_{Q}\left(\left(\nabla_{E_{a}} l\right)(X), E_{a}\right)$. Since $l$ is symmetric,

$$
\begin{equation*}
\sum_{a} E_{a} \cdot l\left(E_{a}\right)=\sum_{a} l\left(E_{a}\right) \cdot E_{a} \tag{5.4}
\end{equation*}
$$

and then

$$
\begin{equation*}
\sum_{a} E_{a} \cdot l\left(E_{a}\right)=-\operatorname{tr} l . \tag{5.5}
\end{equation*}
$$

From (5.4) and (5.5), we have

$$
\begin{equation*}
\sum_{a} E_{a} \cdot\left[l(X), l\left(E_{a}\right)\right] \cdot \Psi=2(\operatorname{tr} l) l(X) \cdot \Psi-2 l^{2}(X) \cdot \Psi . \tag{5.6}
\end{equation*}
$$

From (5.2), (5.3) and (5.6), we have the following proposition.
Proposition 5.2 On the transverse spin foliation, we have

$$
\begin{align*}
\operatorname{Ric}_{\nabla}^{l}(X \otimes \Psi)= & -\frac{1}{2} \rho^{\nabla}(X) \cdot \Psi+\sum_{a}\left(E_{a} \wedge d l\left(X, E_{a}\right)\right) \cdot \Psi-2 l^{2}(X) \cdot \Psi \\
& -\left\{X(\operatorname{tr} l)-\left(d i v_{\nabla} l\right)(X)\right\} \cdot \Psi+2(\operatorname{tr} l) l(X) \cdot \Psi \tag{5.7}
\end{align*}
$$

From (5.7), we have the following facts.
Proposition 5.3 If $M$ admits a non-zero spinor field $\Psi \in \Gamma S(\mathcal{F})$ with $\stackrel{l}{\nabla} \Psi=0$, then $|\Psi|^{2}$ is constant and

$$
\begin{align*}
& \operatorname{grad}_{\nabla}(\operatorname{tr} l)=\operatorname{div}_{\nabla} l,  \tag{5.8}\\
& (\operatorname{tr} l)^{2}=\frac{1}{4} \sigma^{\nabla}+|l|^{2} . \tag{5.9}
\end{align*}
$$

Proof. Since $\stackrel{l}{\nabla}$ is a metric connection, the condition $\stackrel{l}{\nabla} \Psi=0$ imply that $|\Psi|^{2}$ is constant. If $\nabla_{X}^{l} \Psi=0$ for any $X \in \Gamma Q$, then $R i c_{\nabla}^{l}=0$. Hence from (5.7), we have

$$
\begin{gather*}
\left\{X(\operatorname{tr} l)-\left(\operatorname{div_{\nabla }l)(X)\} \Psi }=-\frac{1}{2} \rho^{\nabla}(X) \cdot \Psi+\sum_{a}\left(E_{a} \wedge d l\left(X, E_{a}\right)\right) \cdot \Psi\right.\right. \\
\text { 제⽊ㅜ } 2(\operatorname{tr} l) l(X) \cdot \Psi-2 l^{2}(X) \cdot \Psi
\end{gather*}
$$

Then we have that for any $X \in \Gamma Q$,

$$
\begin{aligned}
& <\left\{X(\operatorname{tr} l)-\left(\operatorname{div}_{\nabla} l\right)(X)\right\} \cdot \Psi, \Psi> \\
& =<\left\{-\frac{1}{2} \rho^{\nabla}(X)+\sum_{a}\left(E_{a} \wedge d l\left(X, E_{a}\right)\right)+2(\operatorname{tr} l) l(X)-2 l^{2}(X)\right\} \cdot \Psi, \Psi>
\end{aligned}
$$

The left-hand side is real, but the right-hand side is pure imaginary because $\sum_{a}<E_{a} \wedge d l\left(X, E_{a}\right) \cdot \Psi, \Psi>$ is pure imaginary. Therefore, both sides are zeros. Hence we get that for any $X \in \Gamma Q$,

$$
\begin{align*}
& X(\operatorname{tr} l)=\left(\operatorname{div}_{\nabla} l\right)(X)  \tag{5.11}\\
& \frac{1}{2} \rho^{\nabla}(X) \cdot \Psi=\left\{\sum_{a}\left(E_{a} \wedge d l\left(X, E_{a}\right)\right)+2(\operatorname{tr} l) l(X)-2 l^{2}(X)\right\} \cdot \Psi . \tag{5.12}
\end{align*}
$$

Hence from (5.11) the equation (5.8) is proved. For the equation (5.9), Clifford multiplication of (5.12) with $E_{b}$ and for $X=E_{b}$, gives

$$
\begin{aligned}
\frac{1}{2} \sum_{b} E_{b} \cdot \rho^{\nabla}\left(E_{b}\right) \cdot \Psi= & \sum_{a, b} E_{b} \cdot\left(E_{a} \wedge d l\left(E_{b}, E_{a}\right)\right) \cdot \Psi \\
& +2(\operatorname{tr} l) \sum_{b} E_{b} \cdot l\left(E_{b}\right) \cdot \Psi-2 \sum_{b} E_{b} \cdot l^{2}\left(E_{b}\right) \cdot \Psi .
\end{aligned}
$$

From Theorem 3.7 and (5.5), we have

$$
\begin{equation*}
-\frac{1}{2} \sigma^{\nabla} \Psi=\sum_{a, b} E_{a} \cdot\left(E_{b} \wedge d l\left(E_{a}, E_{b}\right)\right) \cdot \Psi-2(\operatorname{tr} l)^{2} \Psi+2|l|^{2} \Psi \tag{5.13}
\end{equation*}
$$

On the other hand,

$$
\begin{align*}
& \sum_{a, b} E_{a} \cdot\left(E_{b} \wedge d l\left(E_{a}, E_{b}\right)\right) \cdot \Psi=\left(\sum_{a, b} E_{a} \wedge E_{b} \wedge d l\left(E_{a}, E_{b}\right)\right) \cdot \Psi  \tag{5.14}\\
&-\sum_{\text {i, }} i\left(E_{a}\right)\left(E_{b} \wedge d l\left(E_{a}, E_{b}\right)\right) \cdot \Psi . \\
& \text { 제주대 } \stackrel{a}{\boldsymbol{a}, ~ ㄱ ㅛ ~ ㅈ ㅜ ㅇ ㅇ ㅏ ㅇ ㄷ ㅗ ㅅ ㅓ ㄱ ㅘ ㄴ ~}
\end{align*}
$$

The first term of the right-hand side of (5.14) is zero since $l$ is symmetric, and the last term is the Clifford multiplication of $\Psi$ with a vector field, which gives an imaginary function when taking its scalar product with $\Psi$. Thus we have

$$
-\frac{1}{2} \sigma^{\nabla}|\Psi|^{2}=-2(\operatorname{tr} l)^{2}|\Psi|^{2}+2|l|^{2}|\Psi|^{2}
$$

Hence the proof of (5.9) is completed.
Let $\Psi_{1}$ be the eigenspinor corresponding to the eigenvalue $\lambda_{1}^{2}=$ $\frac{1}{4} \inf _{M}\left(K^{\sigma}+4\left|l_{\Psi_{1}}\right|^{2}\right)$. From (4.12), we have

$$
\begin{equation*}
\stackrel{l_{\Psi_{1}}}{\nabla} \Psi_{1}=0, \quad K^{\sigma}=\text { constant }, \quad\left|l_{\Psi_{1}}\right|=\text { constant } . \tag{5.15}
\end{equation*}
$$

From Proposition 5.3, we know that $\left|\Psi_{1}\right|$ is constant and

$$
\begin{equation*}
\lambda_{1}^{2}=\frac{1}{4} \sigma^{\nabla}+\left|l_{\Psi_{1}}\right|^{2} . \tag{5.16}
\end{equation*}
$$

From (5.16), the transversal scalar curvature $\sigma^{\nabla}$ is constant and we have

$$
\begin{equation*}
\inf |\kappa|^{2}=0 \tag{5.17}
\end{equation*}
$$

Since $\sigma^{\nabla}$ and $K^{\sigma}=\sigma^{\nabla}+|\kappa|^{2}$ are constant, $|\kappa|$ is constant and then $|\kappa|=\inf _{M}|\kappa|=0$. This implies that $\mathcal{F}$ is minimal. Hence we have the following theorem.

Theorem 5.4 Under the same assumption as in Theorem 4.6, if there exists an eigenspinor $\Psi(\neq 0)$ of the basic Dirac operator $D_{b}$ for the eigenvalue $\lambda^{2}=\frac{1}{4} \inf _{M}\left(K^{\sigma}+4\left|l_{\Psi}\right|^{2}\right)$, then $|\Psi|$ is constant and $\mathcal{F}$ is minimal with the constant transversal scalar curvature $\sigma^{\nabla}=4\left(\lambda^{2}-\left|l_{\Psi}\right|^{2}\right)$.

### 5.2 The limiting case II

Next, we study the limiting foliation of (4.35). Similarly, we have that for any $X \in \Gamma Q$,

$$
\begin{align*}
\operatorname{Ric}^{l_{\bar{\nabla}}}(X \otimes \bar{\Phi})= & -\frac{1}{2} \rho^{\bar{\nabla}}(X) \div \bar{\Phi}+\sum\left(\bar{E}_{a} \wedge d l_{\bar{\Phi}}\left(X, \bar{E}_{a}\right)\right) \div \bar{\Phi}-2 l_{\bar{\Phi}}^{2}(X)^{-\bar{\Phi}} \\
& +2\left(\operatorname{tr} l_{\bar{\Phi}}\right) l_{\bar{\Phi}}(X) \div \bar{\Phi}-\left\{X\left(\operatorname{tr} l_{\bar{\Phi}}\right)-\left(d i v_{\bar{\nabla}} l_{\bar{\Phi}}\right)(X)\right\} \cdot \bar{\Phi} \tag{5.18}
\end{align*}
$$

From (5.18), we have that for any $X \in T M$,

$$
\begin{equation*}
\frac{1}{2} \rho^{\bar{\nabla}}(\pi(X))=\sum \bar{E}_{a} \wedge d l_{\bar{\Phi}}\left(\pi(X), \bar{E}_{a}\right)-2 l_{\bar{\Phi}}^{2}(\pi(X))+2\left(\operatorname{tr} l_{\bar{\Phi}}\right) l_{\bar{\Phi}}(\pi(X)) \tag{5.19}
\end{equation*}
$$

Hence we have the following proposition.
Proposition 5.5 If $M$ admits a non-zero spinor $\Psi$ with $\stackrel{{ }_{\bar{\Phi}}}{\nabla} \bar{\Phi}=0$, where
$\Phi=e^{-\frac{q-1}{2} u} \Psi$, then $|\Phi|$ is constant and for any vector field $X$

$$
\begin{align*}
& \nabla_{X} \Psi=\frac{1}{2} \pi(X) \cdot \operatorname{grad}_{\nabla}(u) \cdot \Psi+\frac{q}{2} X(u) \Psi-l_{\Psi}(\pi(X)) \cdot \Psi  \tag{5.20}\\
& X\left(\operatorname{tr} l_{\bar{\Phi}}\right)=\left(\operatorname{div}_{\bar{\nabla}} l_{\bar{\Phi}}\right)(X), \quad\left(\operatorname{tr} l_{\bar{\Phi}}\right)^{2}=\frac{1}{4} \sigma^{\bar{\nabla}}+\left|l_{\bar{\Phi}}\right|^{2} \tag{5.21}
\end{align*}
$$

Proof. Since $\stackrel{l_{\bar{\Phi}}}{\nabla}$ is metrical, $|\Phi|$ is constant. Moreover, ${ }^{l_{\bar{\Phi}}} \bar{\Phi}=0$ is equivalent to

$$
\begin{equation*}
\bar{\nabla}_{X} \bar{\Phi}+l_{\bar{\Phi}}(\pi(X)) \cdot \bar{\Phi}=0 \tag{5.22}
\end{equation*}
$$

From Proposition 3.11 and (5.22), we have that for $\Phi=e^{-\frac{q-1}{2} u} \Psi$

$$
\begin{equation*}
\nabla_{X} \Phi=\frac{1}{2} \pi(X) \cdot \operatorname{grad}_{\nabla}(u) \cdot \Phi+\frac{1}{2} X(u) \Phi-l_{\Psi}(\pi(X)) \cdot \Phi \tag{5.23}
\end{equation*}
$$

which gives (5.20). The proof of (5.21) is similar to the one in Proposition 5.3.

By direct calculation together with (3.18), we obtain the following lemma.

Lemma 5.6 For any vector field $X \in \Gamma Q$ and any isomorphism l, we have

$$
\begin{equation*}
\left(\operatorname{div}_{\bar{\nabla}} l\right)(X)=\left(\operatorname{div}_{\nabla} l\right)(X)+q g_{Q}\left(l(X), \operatorname{grad}_{\nabla}(u)\right)-X(u) \operatorname{tr} l \tag{5.24}
\end{equation*}
$$

where $\bar{\nabla}$ is a Levi-Civita connection with respect to $\bar{g}_{Q}=e^{2 u} g_{Q}$.
On the other hand, we have that for any $\Phi=e^{-\frac{q-1}{2} u} \Psi$ and any $X \in \Gamma Q$

$$
\begin{equation*}
\left(\operatorname{div}_{\nabla} l_{\bar{\Phi}}\right)(X)=e^{-u}\left\{\left(\operatorname{div}_{\nabla} l_{\Psi}\right)(X)-g_{Q}\left(l_{\Psi}(X), \operatorname{grad}_{\nabla}(u)\right)\right\} . \tag{5.25}
\end{equation*}
$$

From (5.24) and (5.25), we have

$$
\begin{align*}
\left(\operatorname{div}_{\bar{\nabla}} l_{\bar{\Phi}}\right)(X)= & e^{-u}\left\{\left(\operatorname{div}_{\nabla} l_{\Psi}\right)(X)+(q-1) g_{Q}\left(l_{\Psi}(X), \operatorname{grad}_{\nabla}(u)\right)\right\}  \tag{5.26}\\
& -e^{-u} X(u) \operatorname{tr} l_{\Psi} .
\end{align*}
$$

Comparing with (5.21) and (5.26), we have the following corollary.

Corollary 5.7 If $M$ admits a non-zero spinor $\Psi$ with ${ }^{l_{\bar{\Phi}}} \bar{\Phi}=0$, where $\Phi=e^{-\frac{q-1}{2} u} \Psi$, then for any $X \in \Gamma Q$

$$
X\left(\operatorname{tr} l_{\Psi}\right)=\left(\operatorname{div}_{\nabla} l_{\Psi}\right)(X)+(q-1) g_{Q}\left(l_{\Psi}(X), \operatorname{grad}_{\nabla}(u)\right) .
$$

Let $D_{b} \Psi=\lambda \Psi$ with $\lambda^{2}=\frac{1}{4}\left(\mu_{1}+\inf _{M}|\kappa|^{2}\right)+\inf _{M}\left|l_{\Psi}\right|^{2}$. From (4.24) and (4.35), we have that for $\Phi=e^{-\frac{q-1}{2} u} \Psi$

$$
\begin{align*}
\mu_{1}+\inf _{M}|\kappa|^{2}+4 \inf _{M}\left|l_{\Psi}\right|^{2} & =\sup _{u} \inf _{M}\left(e^{2 u} K_{\sigma}^{\bar{\nabla}}+4\left|l_{\Psi}\right|^{2}\right) \\
& =\sup _{u \in \mathcal{K}_{u}} \inf _{M}\left(e^{2 u} K_{\sigma}^{\bar{\nabla}}+4\left|l_{\Psi}\right|^{2}\right)  \tag{5.27}\\
& =\inf _{M}\left(e^{2 u} K_{\sigma}^{\bar{\nabla}}+4\left|l_{\Psi}\right|^{2}\right)
\end{align*}
$$

By (4.23),

$$
\begin{equation*}
\stackrel{l_{\Phi}}{\nabla} \bar{\Phi}=0, \quad \inf _{M}\left(e^{2 u} K_{\sigma}^{\nabla}\right)=e^{2 u} K_{\sigma}^{\nabla}, \quad\left|l_{\Psi}\right|=\text { constant. } \tag{5.28}
\end{equation*}
$$

From (5.28), we have

$$
\begin{equation*}
\sup _{u \in \mathcal{K}_{u}} \inf _{M}\left(e^{2 u} K_{\sigma}^{\bar{\nabla}}\right)=\sup _{u \in \mathcal{K}_{u}}\left(e^{2 u} K_{\sigma}^{\bar{\nabla}}\right) \tag{5.29}
\end{equation*}
$$

From (5.29), we have

$$
\begin{equation*}
\sup _{u \in \mathcal{K}_{u}} \inf _{M}\left\{h^{-1} Y_{b} h\right\}+\inf _{M}|\kappa|^{2}=\sup _{u \in \mathcal{K}_{u}}\left\{h^{-1} Y_{b} h\right\}+|\kappa|^{2} \tag{5.30}
\end{equation*}
$$

From (5.27) and (5.30), we have

$$
\begin{equation*}
\mu_{1}=\sup _{u \in \mathcal{K}_{u}} \inf _{M}\left\{h^{-1} Y_{b} h\right\}, \quad|\kappa|=\text { constant. } \tag{5.31}
\end{equation*}
$$

From (4.19) and (5.21), we have

$$
\begin{equation*}
\mu_{1}+\inf _{M}|\kappa|^{2}=e^{2 u} \sigma^{\bar{\nabla}} \tag{5.32}
\end{equation*}
$$

From Lemma 4.13, we have that for $u \in \mathcal{K}_{u}$

$$
\begin{equation*}
e^{2 u} \sigma^{\bar{\nabla}}=\sigma^{\nabla}+2(q-1) \Delta_{B} u+(q-1)(2-q)\left|\operatorname{grad}_{\nabla}(u)\right|^{2} . \tag{5.33}
\end{equation*}
$$

From (4.33), (5.32) and (5.33), we have

$$
\begin{equation*}
\mu_{1}+\inf _{M}|\kappa|^{2}=\sup _{u \in \mathcal{K}_{u}} \inf _{M}\left(e^{2 u} \sigma^{\bar{\nabla}}\right)=\sup _{u \in \mathcal{K}_{u}} \inf _{M}\left\{h^{-1} Y_{b} h\right\}=\mu_{1} . \tag{5.34}
\end{equation*}
$$

From (5.34) we get

$$
\begin{equation*}
\inf _{M}|\kappa|^{2}=0 \tag{5.35}
\end{equation*}
$$

This implies that $|\kappa|=0$, i.e. $\mathcal{F}$ is minimal. Hence we have the following theorem.

Theorem 5.8 Let $\left(M, \tilde{g}_{M}, \mathcal{F}\right)$ be a compact Riemannian manifold with a transverse spin foliation $\mathcal{F}$ of codimension $q \geq 3$ and bundle-like metric $\tilde{g}_{M}$. Assume that an eigenvalue $\lambda$ of $D_{b}$ corresponding to the eigenspinor $\Psi$ satisfies

$$
\lambda^{2}=\frac{1}{4}\left(\mu_{1}+\inf _{M}|\kappa|^{2}\right)+\inf _{M}\left|l_{\Psi}\right|^{2} .
$$

Then $\left|l_{\Psi}\right|$ is constant and the foliation $\mathcal{F}$ is minimal. Moreover

$$
\begin{equation*}
\left(\operatorname{div}_{\nabla} l_{\Psi}\right)(X)=(1-q) g_{Q}\left(l_{\Psi}(X), \operatorname{grad}_{\nabla}(u)\right) \tag{5.36}
\end{equation*}
$$

for any $X \in \Gamma Q$.

Proof. The equation (5.36) is trivial from corollary 5.7.

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## 〈국문초록〉

## 엽층구조를 가지는 리만 다양체에서의 basic Dirac 연산자의 고유치 계산

엽층적 구조를 가지는 리만다양체의 횡단적 구조가 $\operatorname{spin}$ 구조를 가질 때 basic Dirac 연산자의 고유치의 제곱값은 엽층구조 $F$ 의 횡단적 곡률과 평균곡률 및 법속 $Q$ 의 적당한 장기준동형가 $l_{\Psi}$ 의 길이에 종속 되어진다는 것을 보인다. 여기서 $\Psi$ 는 고유 spinor이다. 또한 이 고 유치의 하한은 기존의 알려진 많은 결과들을 모두 유도한다는 것을 보 인다. 더구나, 등식을 만족하는 경우의 엽층의 성질을 조사하여 실제로 엽층들은 모두 극소부분공간임을 알 수 있고, 가장 작은 고유치에 대응 하는 고유 spinor에 의해 결정되어진 자기준동형사상의 길이는 상수라 는 것을 밝혔다.

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